

# Electromagnetically confined plasma target for interaction studies with intense laser fields

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The paper describes a novel application of an electron beam ion trap as a plasma target facility for intense laser–plasma interaction studies. The low density plasma target ( $\sim 10^{13}/\text{cm}^3$ ) is confined in a mobile cryogenic electromagnetic charged particle trap, with the magnetic confinement field of 1–3 T maintained by a superconducting magnet. Ion plasmas for a large variety of ion species and charge states are produced and maintained within the magnetic field and the space-charge of an energetic electron beam in the “electron beam ion trap” (EBIT) geometry. Intense laser beams (optical lasers, X-ray lasers and upcoming “X-ray free electron lasers” (XFEL)) provide strong time varying electromagnetic fields ( $>10^{12}$  V/cm in femto- to nano-second pulses) for interactions with electromagnetically confined neutral/non-neutral plasmas. Experimental scenarios with intense photon fields on ionization/excitation processes, the ionization balance, as well as photon polarization effects and tests with intense lasers that utilize the ion plasma target are outlined. A first test of the plasma target with the PHELIX high intensity laser at GSI is described.

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## 1. Introduction and motivation

The interaction of strong time varying electromagnetic (EM) laser fields with ionized matter are of intrinsic interest to “high energy density physics (HEDP)” ( $>10^{18}$  W/cm<sup>2</sup>) [1,2]. Several fundamental issues of photon–matter interaction physics can be studied with strong EM field pulses interacting with confined and well-defined low density neutral or non-neutral plasmas. Low density plasmas ( $\sim 10^{13}$  ions/cm<sup>3</sup>) will allow to isolate and identify particular excitation processes (e.g. excitation/recombination) and to study them as a function of the laser interaction. This in turn provides benchmark data for the more complex case of dense matter under extreme conditions.

Ionization/excitation are the primary mechanism by which photons deliver heat to a sample and create high density plasmas. High intensity laser interaction with plasmas

produces electrons at high energies (keV to MeV) via various overlapping processes (e.g. plasma instabilities), which effect the ionization balance and “equation of state” (EOS). The interpretation of the EOS for hot dense plasmas requires data of the ionization stages to determine the specific heat contributions. A question of fundamental interest is, for example, to what extent a low density, equilibrium or non-equilibrium photo-ionized plasma consisting of multiple ionized atoms can be created and be probed in the density-independent, radiation-dominated regime (e.g. in relevance to neutron-star accretion disk physics).

In short, pulsed laser interactions, multiple ionization of atoms, using intense optical fs lasers, has been studied experimentally and theoretically resulting in mostly qualitative data [3]. Many surprises were found and models are revised continuously and range from a sequential ionization [4], i.e. 2+ comes from 1+ and 3+ via 2+, etc., to the tunneling ionization model and above threshold ionization, to a re-scattering model of the ionized electron that causes multiple ionization [5]. In these cases, the strong electrical field modulates the Coulomb potential and subsequent ionization processes are responsible for the degree of ionization [1]. With optical lasers, the competition between sequential and non-sequential processes has been investigated and is interpreted on the basis of tunneling and many-body S-matrix scattering theory as well as applying e.g. the Lotz formula. The presence of a high ac electric field requires the field to be treated as a strong perturbation; the Stark effect, e.g. in atoms (resonant), can be described by coupled light-matter states (“dressed” states). In non-resonant cases, non-perturbative phenomena occur, such as above-threshold ionization, higher-order harmonic generation and the laser assisted photoelectric effect. The non-resonant, non-perturbative effects occur when the applied field is at the transition between the classical and quantum regime (limit of single photon effect). A measure for the field strength that defines this regime is the ponderomotive potential,  $U_p$ . It represents the averaged “quiver” kinetic energy of an electron in an ac electric field.  $U_p$  is given as

$$U_p = (e^2 A_0^2) / (4mc^2) = \left[ (2\pi e^2) / (mc) \right] (I / \omega^2)$$
 with I being the intensity of the electric field at the frequency  $\omega$  and  $A_0^2$  the amplitude of the electric field vector potential. If  $U_p$  is comparable to e.g. the ionization energy ( $E_{IP}$ ) non-perturbative effects appear (see

$\gamma = \sqrt{E_{IP} / U_p} < 1$ , the Keldysh parameter [6]). It is noted that the “quiver” energies can reach  $\sim 3$  MeV, leading to relativistic degenerated electrons. The majority of data have been obtained for neutral targets and clusters [7], where the latter shows enhanced ionization probabilities. For clusters, an important phenomenon is the strongly increased energy absorption compared to the single-atom case. The mechanisms, depending on the cluster size, are the enhanced ionization for small clusters, resonant coupling of the laser frequency to a collectively oscillating electron cloud in intermediate clusters, and the effective energy absorption by plasma oscillations resonantly coupled to the laser frequency in large clusters. Laser interactions with clusters or micro-plasmas can be viewed to some extent as an intermediate situation between the single-atom case and complex plasmas. Very little work has been done so far with optical high intensity lasers on the ionization of ions and in particular highly charged ions as they are present in HEDP cases. So far, only singly charged ions were investigated with strong laser fields. Recently, data on singly and low charged ions have been published, showing a strong suppression of non-sequential processes. Evidence for the re-collision model has been

found from the suppression of non-sequential processes using circularly polarized fields and from analyzing the recoil ion momentum distribution [8–13]. In the re-collision process, the electrons may occupy bound states of the parent atom and then release energy in form of short coherent EUV or X-ray bursts. The re-colliding electrons are “wiggling” in the Coulomb field of the parent atoms, generating higher harmonics with potentially very high harmonic numbers with multiply-charged ions depending on the laser determined re-scattering process and the ion species.

The paper describes the novel plasma target facility for the experimental investigation of interactions of strong electromagnetic fields with an electromagnetically confined neutral/non-neutral plasma at adjustable densities ( $<10^{13}/\text{cm}^3$ ). The laser used for the first test experiments described here is the short pulsed preamplifier of the PHELIX (peta-watt) laser [14] at GSI-Darmstadt. 2. The plasma target facility A plasma target device based on the design of an “electron beam ion trap (EBIT)” has the advantage that one can breed charge states and study photon field induced multiple ionization/excitation starting from different initial charge states. The studies can be performed by observing photon emission from the plasma in situ or by analyzing the extracted ion charge state distribution. Electron beam ion trap style sources have been used for years to produce low emittance, low kinetic energy beams of highly charged ions [15–17] for surface physics studies, for re-trapping and laser cooling of highly charged ions [18,19] and for in situ X-ray spectroscopy [20,21]. The specially designed electron beam ion trap (Fig. 1) as an ion or ion/electron low density plasma ( $<10^{12}/\text{cm}^3$ ) target facility resembles an open cylinder Penning trap geometry, where an electron beam passes through three axial drift tubes as trap electrodes inside a cryogenic super-conducting 3 T magnet. The plasma is produced and confined in the center drift tube. The ions are confined radially by the space-charge of the compressed, high density ( $10^{12}/\text{cm}^3$ ) electron beam and magnetic field while axial confinement is provided by appropriate biasing the trap drift tubes; the trap biases are superimposed on a high voltage potential which determines the electron beam impact energy. The electron beam energy, intensity and the ionization time determine the achievable charge state distribution via successive electron impact ionization. The target facility provides selectively narrow charge state distributions of highly charged ions up to about  $\text{U}^{80+}$ . The base vacuum in the plasma region is maintained at  $10^{-11}$  Torr. The ion plasma target is equipped with especially designed windows and apertures to allow the high power laser beams to interact with the plasma without distorting the vacuum or causing background effects. The plasma target device has an innovative design, using new cryogenic vacuum technology, which allows the trap core to remain at the nominal super-conducting operation temperature without relying on liquid gas cooling [22,23] and allows easy access as well as transfer to different target locations. During testing and performance verification [22], highly charged Ar- and Xe ions have been produced in a commercially available device temporarily used for testing at the Lawrence Berkeley Laboratory [19]. In that device the ions were produced using up to a 15 keV electron beam energy at a beam current of 50–100 mA, a confinement time of 500 ms and a saw-tooth extraction pulse with an extraction base width of 100 ms. The extracted ions are focused, momentum analyzed in a  $90^\circ$  analyzing magnet and detected with a multi channel plate (MCP) detector and Faraday cup. The measured ion yields for  $\text{Ar}^{18+}$  and  $\text{Xe}^{45+}$  ions are  $10^7/\text{s}$  in a less than 1 mm spot 2.5 m from the trap; the emittance of the ion beams is  $0.3 \pi$  mm mrad (see e.g. [24]).

Different experimental scenarios can be arranged for laser plasma interaction studies. Laser injection can be accomplished by focusing the laser through side windows onto the target. A coaxial laser arrangement with the plasma, along the magnetic field axis of the trap can be arranged as well. Higher density target plasmas are generally obtainable by increasing the electron beam current and magnetic field. It is noted that these increases will significantly raise the plasma ion temperature prior to the laser interaction. The open trap geometry and electron beam space charge trapping can also be used to produce plasma targets from radially inserted solid microprobes or from a gas “puff” target.

### 3. Experiment scenarios and first tests

A scenario for the laser/plasma interaction studies has been tested in a first experiment at the GSI PHELIX laser. The EBIT that was used to produce the plasma target, has been provided temporarily by the University of Stockholm (Atomic Physics Department). For the tests a high intensity laser beam ( $10^{14}$  W/cm<sup>2</sup>) has been focused onto the confined electron beam/ion plasma target ( $e + \text{Xe}^{9+}$ ) through a side window (Fig. 1). A ray-tracing simulation of the laser injection set-up was performed to analyze the possibility to change the circular point focus (e.g. 30  $\mu\text{m}$ ) to a well-defined elliptical focus, matching the target geometry. In the first test experiments, low energy photon emission, X-ray emission and extracted ions have been detected with and without laser interaction. A solar blind photomultiplier was mounted behind a quartz window on top of the EBIT chamber to detect the emission of low energy photons perpendicular to the interaction region. The broad energy range of the photo multiplier of <200 eV is determined by the window’s transmission and the multiplier detection efficiency. A small solid angle X-ray Si-detector (AMTEK) viewed the same interaction region as the photomultiplier through a Be window at another side port on the EBIT. Highly charged ions were extracted from the plasma in a pulsed mode by ramping the center drift tube potential with and without laser interaction. The ions were momentum analyzed in a 30° bending magnet and detected with a position sensitive channel plate detector. Extracted ions were also analyzed via their emission of radiative de-excitation (“hollow” atom) radiation following the surface interaction at a thin Be window on a straight through beam line port in the analyzing magnet; these X-rays were detected with a Ge detector. The photon emission and ion extraction detection has been gated to the laser pulse timing and the data acquisition was done in event mode. The laser focus of e.g. 30  $\mu\text{m}$  diameter matches the ion volume in EBIT of nominal 60  $\mu\text{m}$  diameter at 3 T. Laser beams of 0.4 ps pulse length of typical  $10^{14}$  W/cm<sup>2</sup> power density from the PHELIX preamplifier (0.01 Hz) and of low power density ( $10^{12}$  W/cm<sup>2</sup>) from the PHELIX frontend at 10 Hz have been used. These first tests demonstrated the capability to focus a high power laser beam onto an electron beam/ion plasma target without distortion of the sensitive plasma confinement conditions (electromagnetic field, 4 K temperature, vacuum <10<sup>11</sup> Torr). Several tests have been performed to ensure that the laser intensity is focused onto the confined electron ion plasma target without scattering at the entrance slits at the center drift tube or any other aperture in the laser path. The tests showed, that it will be possible to observe (e.g. in coincidence) the X-ray emission and ion extraction from

the confined ion plasma. The application of the different laser conditions will explore to what extent the strong time varying electromagnetic field pulses of the laser affect the ionization stages and their temporal development in the ion plasma. First evidence for significant effects on the photon emission from the plasma was observed which indicates the possibility to identify “signature” processes for this type of laser/plasma interaction. The relevant laser parameters are given in Table 1 together with the ion/electron plasma target parameters.

We summarize some of the evidence for effects from the first laser/plasma interaction experiments as follows:

(1) Photon detection in photomultiplier on top viewport:

Large photon (<200 eV) bursts following the high power (HP) laser beam interaction at  $>10^{14}$  W/cm<sup>2</sup> with an electron beam/Xe<sup>9+</sup> ion plasma, are observed. No photon emission was observed when the electron beam was turned off. The HP laser has a repetition rate of 3 min. In between the HP laser pulses, a low power (LP) laser at  $10^{12}$ W/cm<sup>2</sup> with a repetition rate of 10 Hz was turned on and smaller (factor 10) photon “bursts” were observed. These intensities appeared to decay exponentially over 3 min (HP laser repetition rate). Each laser pulse is followed by an ion extraction cycle of 150 ms and “reloading” (300 ms) of the trap, therefore, the long decay time is indicating an interaction with the electron population in the trap. This observation might be explained by Thomson (up-) scattering of laser photons including excitation from the electron beam electrons into high n states in the 3 T field (Landau states?). The decaying photon intensities may reflect the decreasing number of excited electrons in the laser interaction region inside the trap volume.

(2) Highly charged Xe ions have been extracted and used to bombard a thin Be window at the end of the straight beam line section in the analyzing magnet chamber. The observation of radiative de-excitation X-rays confirms the extraction of highly charged Xe ions centered around 44+ in agreement with earlier results. No effect from the laser interaction (e.g. energy shift due to change in the charge state distribution of highly charged Xe ions) could be seen due to low counting statistics. The charge state distribution could vary due to effects on the ion heating and an effective change of the trapping potential caused by the laser interaction.

(3) For the momentum analyzed ions, the overall yield of extracted ions seemed to be increased through the presence of the laser interaction (about 20%) which could be due to a temporary change of the effective trapping potential. This effect could, if verified in future experiments, be interesting for creating short pulsed HCI extraction with subsequent short pulsed X-ray emission following radiative de-excitation in a solid surface target.

#### **4. Future experiments**

At XFEL facilities (e.g. [25,26]), higher photon energies, intensities and repetition rates are available, which enable the conductance of a new class of photon/confined multiple- charged ion plasma interaction studies. It is possible to investigate the interaction of high intensity, short wavelength, up to 8 keV (0.15–1.5 nm) photon beams

(LCLS) with low density electromagnetically confined neutral/non-neutral multiple-charged ion/electron plasmas. Strong field and X-ray induced/assisted atomic multiple ionization/ excitation processes, could possibly be investigated. For example, it is expected to observe effects from the strong electromagnetic laser fields on the ion charge state distribution in the plasma with and without the presence of the electron beam. The effects will be detected by momentum analyzing the extracted plasma ions and by measuring the X-ray emission from the interaction region. Some expected XFEL (LCLS) parameters are given in Table 1 for comparison with the (PHELIX) laser parameters.

Another class of experiments could utilize the extraction of HCI beams and their interaction with surfaces. Here the ions would act as ‘‘pump’’, for example. Large cluster emission has been reported from HCI impact and subsequent laser cluster interaction studies could be conducted on uniquely formed clusters such as  $(\text{SiO}_2)_{22}$  and  $(\text{UO}_2)_{14}$ . Other studies could utilize the formation of dense, rather hot micro-plasmas at the surface impact sites of HCIs. A probing laser beam focused onto the micro-plasma would allow to study correlation and inverse bremsstrahlung effects in such (strongly coupled, ‘‘warm dense matter’’) plasmas.

## 5. Conclusion

We demonstrated the feasibility of using an electron beam ion trap as a plasma target facility for intense laser beam interaction studies. The laser experiments are aimed at gaining understanding of the effects of intense photon fields on ionization/excitation processes, the ionization balance, as well as photon polarization effects on confined charged particle plasmas. The low density plasma target ( $\sim 10^{13}/\text{cm}^3$ ) is confined in a cryogenic electromagnetic trap, with the magnetic confinement field of 1–3 T maintained by a superconducting magnet. First experimental scenarios have been demonstrated and various tests have been performed, which confirmed functioning of all the plasma confinement conditions (e.g. trapping potential) during the intense laser interaction. The tests also showed that it will be possible to observe (simultaneously or in coincidence) the X-ray emission and ion extraction from the confined ion plasma during the laser interaction.

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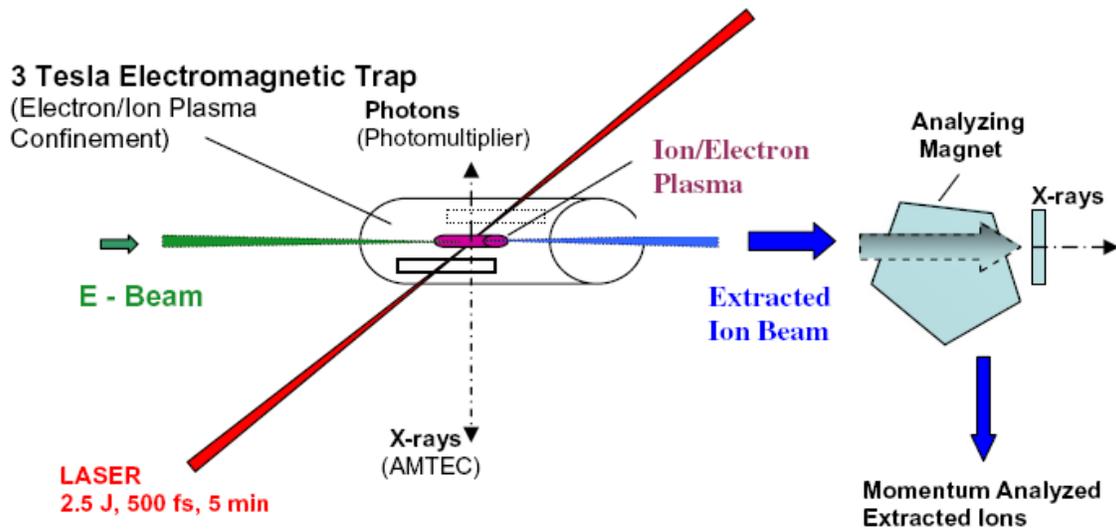


Fig. 1. Schematic of experimental set-up.

Table 1  
Relevant plasma target, laser and expected XFEL (LCLS) parameters

<i>Electron/ion plasma target</i>	
Magnetic field	3 T
Electron beam energy	30 keV
Electron beam current	150 mA
Electron beam radius	37 $\mu\text{m}$
(Trapped ion volume)	
Central current density	4 kA/cm <sup>2</sup>
Trap length	2 cm
Electron density	10 <sup>11</sup> /cm <sup>3</sup>
Ion density	10 <sup>11</sup> /cm <sup>3</sup>
<i>PHELIX laser preamplifier</i>	
Wave length	1054 nm
Energy	0.2–5 J
Focus	30 $\mu\text{m}$
Pulse duration	500 fs (long pulse 10 ns)
Repetition rate	3 min
E-field	1.7 $\times 10^{10}$ V/cm
B-field:	5800 T
<i>LCLS (linac coherent light source)</i>	
Wavelength	0.15–0.15 nm
Peak (average) brightness	10 <sup>33</sup> (average 10 <sup>20</sup> ) photons s <sup>-1</sup> mm <sup>-2</sup> mr <sup>-2</sup> (0.1% bandwidth) <sup>-1</sup>
Pulse duration	<230 fs
Repetition rate	120 Hz