

Strontium and barium iodide high light yield scintillators

Nerine J. Cherepy,^{1,a)} Giulia Hull,¹ Alexander D. Drobshoff,¹ Stephen A. Payne,¹ Edgar van Loef,² Cody M. Wilson,² Kanai S. Shah,² Utpal N. Roy,³ Arnold Burger,³ Lynn A. Boatner,⁴ Woon-Seng Choong,⁵ and William W. Moses⁵

¹Lawrence Livermore National Laboratory, Livermore, California 94550, USA

²Radiation Monitoring Devices, Boston, Massachusetts 02134, USA

³Center for Physics and Chemistry of Materials, Fisk University, Nashville, Tennessee 37208, USA

⁴Oak Ridge National Laboratory, Oak Ridge, Tennessee 37831, USA

⁵Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA

Europium-doped strontium and barium iodide are found to be readily growable by the Bridgman method and to produce high scintillation light yields. $\text{SrI}_2(\text{Eu})$ emits into the Eu^{2+} band, centered at 435 nm, with a decay time of 1.2 μs and a light yield of $\sim 90\,000$ photons/MeV. It offers energy resolution better than 4% full width at half maximum at 662 keV, and exhibits excellent light yield proportionality. $\text{BaI}_2(\text{Eu})$ produces $>30\,000$ photons/MeV into the Eu^{2+} band at 420 nm ($<1\ \mu\text{s}$ decay). An additional broad impurity-mediated recombination band is present at 550 nm ($>3\ \mu\text{s}$ decay), unless high-purity feedstock is used.

Detection sensitivity for weak gamma ray sources and rapid unambiguous isotope identification are principally dependent on energy resolution, and are also enhanced by high effective atomic number of the detector material. The inorganic scintillator currently providing the highest energy resolution is $\text{LaBr}_3(\text{Ce})$, $\sim 2.6\%$ at 662 keV,¹⁻³ but it is highly hygroscopic, possesses intrinsic radioactivity due to the presence of primordial ^{138}La , and its crystal growth is still challenging. Strontium and barium iodide doped with europium are readily growable orthorhombic crystals that offer low phonon frequencies, moderate density, $\rho=4.6$ and $5.1\ \text{g}/\text{cm}^3$, respectively, equivalent or higher light yields than $\text{LaBr}_3(\text{Ce})$ and no intrinsic radioactivity.

Reports of alkaline earth halide scintillation originate with Hofstadter *et al.* on calcium iodide in the 1960s.⁴ Calcium iodide exhibits light yields of $\sim 100\,000$ photons/MeV and has been activated with Tl^+ and Eu^{2+} ; however, it is nearly impossible to grow substantial CaI_2 crystals due to its platelet growth habit. While Hofstadter *et al.* patented the $\text{SrI}_2(\text{Eu})$ crystal in 1968,⁵ no isotope-identifying devices based on this material were ever reported. A report on cathodoluminescence from Ca, Sr, and Li halides described efficient Eu^{2+} activation and moderate hygroscopicity.⁶ A few studies of scintillation from undoped BaI_2 have appeared.^{7,8} The efforts of Selling *et al.* to observe scintillation from Eu^{2+} doped BaI_2 reported a light yield of 2000 photons/MeV.⁹ In recent years, this class of materials has been largely ignored for scintillation applications.

Strontium iodide was grown at Radiation Monitoring Devices (RMD) while barium iodide crystals were grown at Fisk University, Oak Ridge National Laboratory and RMD; all in quartz crucibles using the Bridgman method. The melting points of SrI_2 and BaI_2 are 515 and 711 $^\circ\text{C}$, respectively; both possess orthorhombic symmetry^{9,10} (while calcium iodide is hexagonal⁴). All crystals described in this letter were doped with 0.5 mole % europium and were several cubic centimeters per boule, then cut into $\sim 1\ \text{cm}^3$ pieces for evalu-

ation. Barium iodide as-supplied powder, 99.995% pure ultradry (Alfa Aesar) was yellowish in color (thought to be due to oxide or oxyiodide contamination). Crystals grown directly from as-supplied powders retained a dark coloration (referred to henceforth as “first crystal”). Zone refining rendered the starting powders colorless, and the resulting pure powders were used to grow several crystals (referred to as “second crystal,” although several were grown following this procedure). Finally, an ultrapurification method was used at RMD to grow a $\text{BaI}_2(\text{Eu})$ crystal, referred to as “third crystal.”

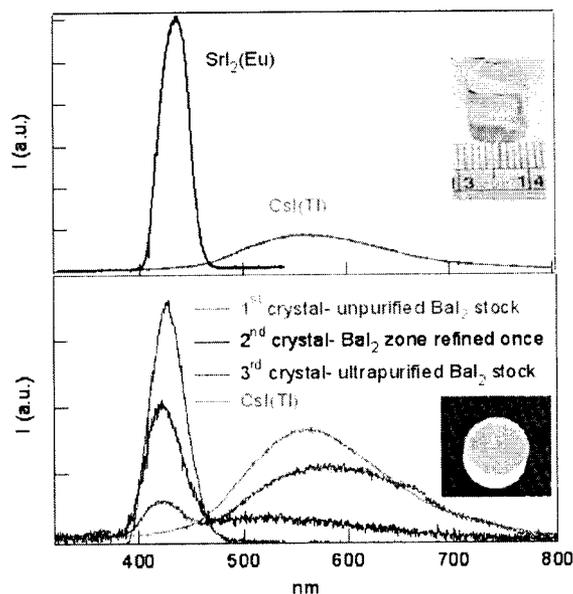


FIG. 1. (Color online) (Top frame) Beta-excited radioluminescence spectra acquired of $\text{SrI}_2(\text{Eu})$; (inset) $\text{SrI}_2(\text{Eu})$ crystal grown at RMD under 254 nm excitation. (Lower frame) Beta-excited radioluminescence from three $\text{BaI}_2(\text{Eu})$ crystals; (inset) $\text{BaI}_2(\text{Eu})$ crystal from Fisk University under 254 nm excitation. For $\text{BaI}_2(\text{Eu})$, the first crystal was grown from as-received powders (red), the second crystal was grown from powders that had been zone refined once (black), and the third crystal was grown under ultrapure conditions. Light yields are plotted with that of $\text{CsI}(\text{Tl})$ (green), and may be compared in an absolute sense.

^{a)}Tel.: 925-424-3492. FAX: 925-423-8772. Electronic mail: cherepy1@llnl.gov.

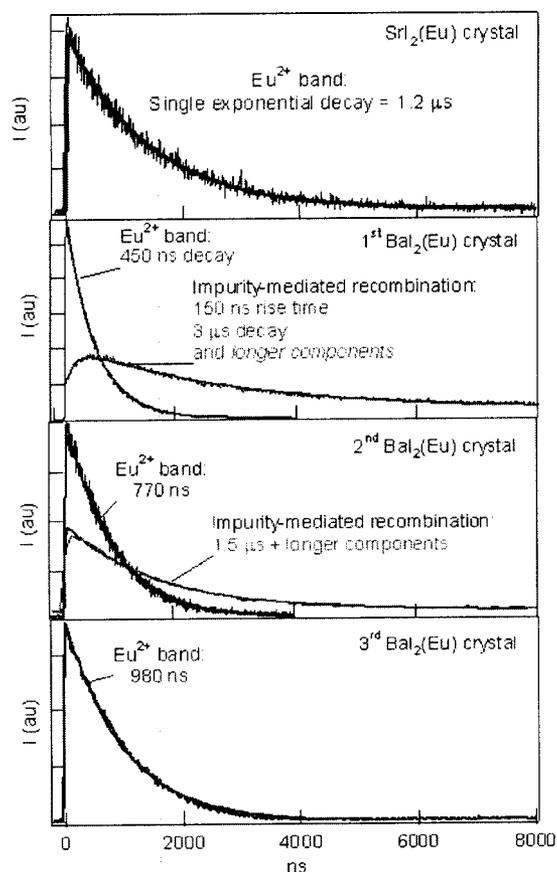


FIG. 2. (Color online) Time-resolved luminescence decays acquired by excitation with 30 ns laser pulses at 266 nm. $\text{BaI}_2(\text{Eu})$ emission recorded at 420 nm for Eu^{2+} emission and at 600 nm for “impurity recombination” band.

Radioluminescence spectra were acquired using a $^{90}\text{Sr}/^{90}\text{Y}$ source (average beta energy ~ 1 MeV) to provide a spectrum expected to be essentially equivalent to that produced by gamma excitation. Radioluminescence spectra were collected with a spectrograph coupled to a thermoelectrically cooled camera and corrected for spectral sensitivity. The beta-excited luminescence of $\text{SrI}_2(0.5\% \text{Eu})$ compared to that of a standard scintillator crystal, $\text{CsI}(\text{Tl})$, is shown in Fig. 1 (top frame), along with a $\text{SrI}_2(\text{Eu})$ crystal grown at RMD (inset). It possesses a single band centered at 435 nm, assigned to the $\text{Eu}^{2+} d \rightarrow f$ transition, and an integrated light yield of 93 000 photons/MeV. Figure 1 (bottom frame) shows beta-excited luminescence spectra of three $\text{BaI}_2(\text{Eu})$ crystals compared to a $\text{CsI}(\text{Tl})$ standard crystal. The Eu^{2+} luminescence at ~ 420 nm is enhanced in the second $\text{BaI}_2(\text{Eu})$ crystal, while the ~ 550 nm band is reduced, and for the third crystal, the ~ 550 nm band is entirely absent. It is notable that the overall light yield is highest for the first crystal; its integral light yield (including both the 420 nm and the 550 nm bands) is 60 000 photons/MeV. The weak band at 550 nm may be assigned to an impurity-mediated recombination transition.

Decay lifetimes were acquired using a flashlamp-pumped Nd:yttrium aluminum garnet laser at 266 nm with 20 ns pulses. Luminescence is collected with a monochromator coupled to an R928 Hamamatsu PMT and readout by an oscilloscope. In $\text{SrI}_2(\text{Eu})$, the Eu^{2+} band exhibits a $1.2 \mu\text{s}$ decay (Fig. 2, top). Figure 2, second from top, shows that for

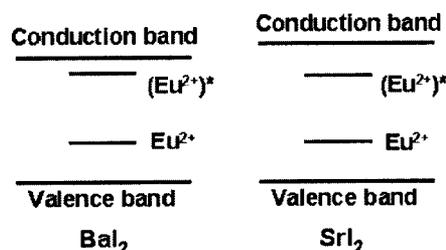


FIG. 3. (Color online) Postulated energy level diagram for the $\text{BaI}_2(\text{Eu})$ and $\text{SrI}_2(\text{Eu})$ systems.

the first $\text{BaI}_2(\text{Eu})$ crystal, the Eu^{2+} decay is ~ 450 ns while the impurity-mediated luminescence is slower, containing components that are longer than $8 \mu\text{s}$. It is interesting that a component of the impurity-mediated recombination is prompt; an additional component forms by depopulating the Eu^{2+} excited state (possibly electrons trapped initially at Eu^{2+} thermally detrapp to the conduction band), as revealed by a rise-time component. For the second $\text{BaI}_2(\text{Eu})$, the Eu^{2+} decay is ~ 770 ns, as shown in Fig. 3 (third from top), effectively lengthened due to the reduction of detrapping and excitation transfer to the impurity-mediated recombination route. The third $\text{BaI}_2(\text{Eu})$ crystal exhibits an even slightly longer decay, ~ 980 ns, as impurity-mediated recombination has been eliminated. A diagram indicating a proposed energy level diagram for $\text{SrI}_2(\text{Eu})$ and $\text{BaI}_2(\text{Eu})$ is shown in Fig. 3.

Gamma ray spectra were acquired using a Hamamatsu R980 alkali photomultiplier tube (PMT) (spectral sensitivity in 380–420 nm range is nearly constant $\sim 30\%$). Crystals were centered on the entrance window, optically coupled to the PMT with mineral oil and wrapped with several layers of Teflon tape. The PMT signals were shaped with a Tennelec TC 244 spectroscopy amplifier [$4 \mu\text{s}$ shaping time for $\text{SrI}_2(\text{Eu})$ and $8 \mu\text{s}$ for $\text{BaI}_2(\text{Eu})$], then recorded with an Amptek MCA8000-A multichannel analyzer. The total gamma absorption peaks (“photopeaks”) were fit to a Gaussian to evaluate the peak position and FWHM to estimate the scintillation light yield and the energy resolution, respectively. Gamma light yields are determined by direct comparison of the photopeak position for $\text{SrI}_2(\text{Eu})$ (channel No. 376;

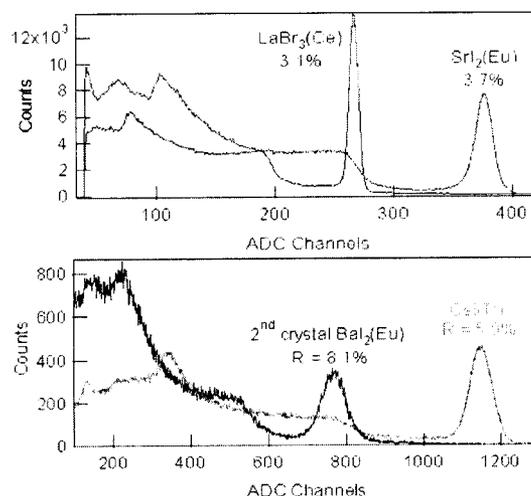


FIG. 4. (Color online) Pulse-height spectra acquired of a $\text{SrI}_2(\text{Eu})$ crystal grown at RMD (upper frame), yielding an energy resolution of 3.7% at 662 keV, and of the second crystal of $\text{BaI}_2(\text{Eu})$ (lower frame), yielding an energy resolution at 662 keV of 8.1%.

TABLE I. Beta-excited light yields are acquired with a silicon CCD camera that integrates the luminescence spectrum over 30 s; spectral sensitivity corrected for CCD and grating. Gamma-excited light yields are determined via pulse-height spectra, corrected for spectral sensitivity, and acquired at 662 keV using a Hamamatsu R980 PMT (4–8 μ s shaping time).

	Gamma light yield (photons/MeV)	Beta light yield (photons/MeV)	Energy resolution (662 keV)
CsI(Tl)	65 000 ^a	65 000 ^b	6.2%
LaBr ₃ (Ce)	60 000 ^a	60 000 ^b	3.2%
SrI ₂ (Eu)	85 000	93 000	3.7%
(First crystal) BaI ₂ (Eu)	14 000	60 000 ^c	8.1%
(Second crystal) BaI ₂ (Eu)	10 000	35 000 ^c	8.1%
(Third crystal) BaI ₂ (Eu)	35 000	40 000	>10%

^aAs quoted by crystal supplier.

^bAssumed to be the same for betas and gammas.

^cIncludes both 420 and 550 nm luminescence bands.

85 000 photons/MeV) and LaBr₃(Ce) (channel No. 266; 60 000 photons/MeV). Figure 4 (top) shows the pulse-height spectra acquired using the 662 keV gamma from ¹³⁷Cs for SrI₂(Eu) and LaBr₃(Ce) under the same conditions. Energy resolution at 662 keV of <4% and light yield significantly superior to that of LaBr₃(Ce) are reproducibly measured for SrI₂(Eu). Subsequent SrI₂(Eu) crystals grown by ORNL, with Eu doping of 4% and 6%, exhibit gamma light yields in excess of 100 000 photons/MeV, and similar energy resolution. In Fig. 4 (bottom), the pulse-height spectrum of the first BaI₂(Eu) crystal is reported, in comparison to CsI(Tl). Energy resolution at 662 keV for the second BaI₂(Eu) crystal was measured to be 8%. While the third BaI₂(Eu) crystal exhibited a gamma light yield of 35 000 photons/MeV, its energy resolution was poor, possibly due to nonuniformity of doping and suboptimal optical quality. Table I summarizes light yield and resolution results. For BaI₂(Eu), the discrepancy between the beta and gamma integral light yields is due to delayed luminescence components. Figure 5 shows the energy resolution as a function of gamma ray energy for SrI₂(Eu) and LaBr₃(Ce) using Ba-133, Am-241, Co-57, Na-22, Co-60, and Cs-137 sources along with a fit to Poisson statistics.

The scintillation light yield nonproportionality characterization instrument (SLYNCI), a unique facility for measuring the so-called nonproportionality of scintillator materials, is described in Ref. 11. Figure 6 shows the relative light yield as a function of electron energy for SrI₂(Eu), compared to that of NaI(Tl) and LaBr₃(Ce).^{12,13} The light yield propor-

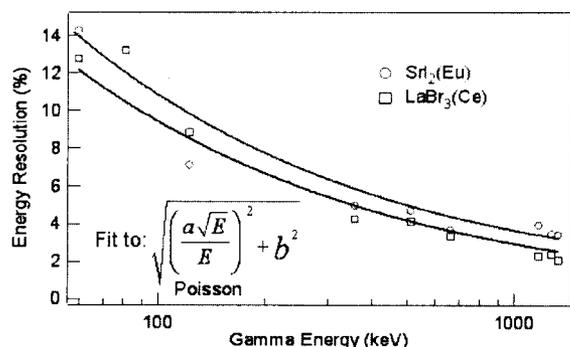


FIG. 5. (Color online) Pulse-height spectra provide the energy resolution as a function of gamma ray energy. Energy resolution is comparable between LaBr₃(Ce) and SrI₂(Eu) for all energies.

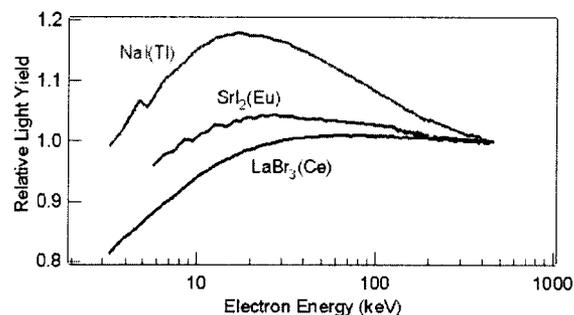


FIG. 6. (Color online) Relative light yields as a function of electron energy acquired using the SLYNCI reveal a proportional response for SrI₂(Eu), in comparison to both LaBr₃(Ce) and NaI(Tl).

tionality is excellent for SrI₂(Eu), and thus the contribution to energy resolution from nonproportionality may be small for SrI₂(Eu). This may be due to extremely efficient capture of excitation on Eu²⁺ sites, more independent of excitation density than for other scintillators.^{12,13} Future experiments will be needed to verify this expectation.

In summary, strontium iodide is a readily “growable” crystal that activates efficiently with Eu²⁺, yielding >80 000 photons/MeV and demonstrating <4% energy resolution at 662 keV. Its energy resolution and light yield proportionality surpass NaI(Tl) and approach LaBr₃(Ce). These initial results are very promising considering that the crystal uniformity, light collection, and readout can still be optimized. We are exploring improvements in crystal growth technique and higher Eu²⁺ doping in order to evaluate the ultimate performance of SrI₂(Eu) and BaI₂(Eu) for high energy resolution gamma ray spectroscopy.

This work was supported by the Domestic Nuclear Detection Office in the Department of Homeland Security and by the NNSA, Office of Nonproliferation Research and Development (NA-22) of the U.S. Department of Energy under Contract No. DE-AC03-76SF00098, and was performed under the auspices of the U.S. DOE by Lawrence Livermore National Laboratory under Contract No. DE-AC52-07NA27344. Oak Ridge National Laboratory is managed for the U.S. DOE by UT-Battelle under Contract No. DE-AC05-00OR22725.

¹E. V. D. van Loef, P. Dorenbos, C. W. E. van Eijk, K. Kramer, and H. H. Güdel, *Appl. Phys. Lett.* **79**, 1573 (2001).

²K. S. Shah, J. Glodo, M. Klugerman, W. M. Higgins, T. Gupta, and P. Wong, *IEEE Trans. Nucl. Sci.* **51**, 2395 (2004).

³C. Fiorini, A. Gola, M. Zanchi, A. Longoni, P. Lechner, H. Soltau, and L. Strüder, *IEEE Trans. Nucl. Sci.* **53**, 2392 (2006).

⁴R. Hofstadter, E. W. O’Dell, and S. T. Schmidt, *IEEE Trans. Nucl. Sci.* **NS-11**, 12 (1964).

⁵R. Hofstadter, U.S. Patent No. 3,373,279 (2 March 1968).

⁶W. Lehmann, *J. Electrochem. Soc.* **122**, 748 (1975).

⁷S. Kubota, J. Ruan (Gen), M. Itoh, S. Hashimoto, and S. Sakuragi, *Nucl. Instrum. Methods Phys. Res. A* **289**, 253 (1990).

⁸P. A. Rodnyi, *Radiat. Meas.* **33** 605 (2001).

⁹J. Selling, M. D. Birowosuto, P. Dorenbos, and S. Schweizer, *J. Appl. Phys.* **101**, 034901 (2007).

¹⁰V. Barnighausen and N. Schultz, *Acta Crystallogr., Sect. B: Struct. Crystallogr. Cryst. Chem.* **B25**, 1104 (1969).

¹¹W.-S. Choong, K. M. Vetter, W. W. Moses, S. A. Payne, N. J. Cherepy, J. D. Valentine, and G. Hull, *IEEE Trans. Nucl. Sci.* (unpublished).

¹²W.-S. Choong, G. Hull, W. W. Moses, S. A. Payne, N. J. Cherepy, J. D. Valentine, and B. W. Reutter, *IEEE Trans. Nucl. Sci.* (unpublished).

¹³G. Hull, W.-S. Choong, G. Bizarri, W. W. Moses, J. D. Valentine, S. Payne, N. Cherepy, and B. W. Reutter, *IEEE Trans. Nucl. Sci.* (unpublished).