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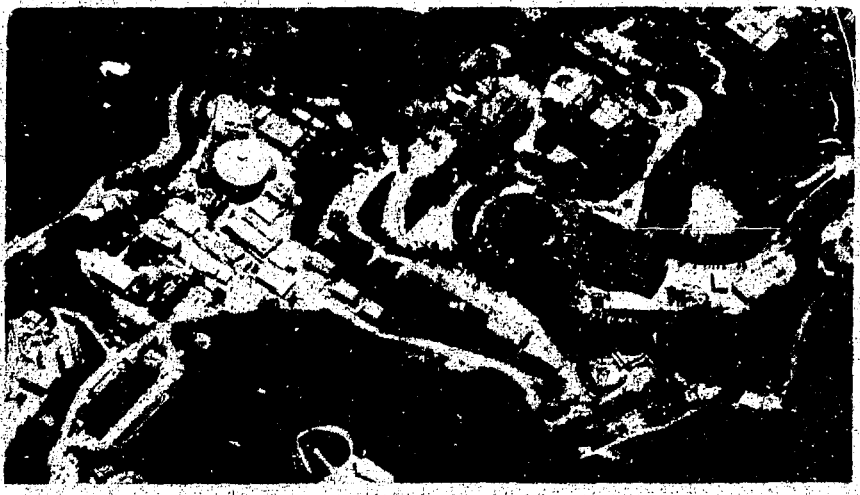
Lawrence Berkeley Laboratory
UNIVERSITY OF CALIFORNIA

Physics, Computer Science & Mathematics Division

STUDY OF RARE PROCESSES INDUCED BY 209-GeV MUONS

Wesley Harold Smith
(Ph.D. thesis)

May 1981



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Nos. W-7405-ENG-48, DE-AC02-76ER03072, and EY-76-C-02-3000

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Contents

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A Study of Rare Processes Induced by 209-Gev Muons

Wesley H. Smith

ABSTRACT

Analysis of dimuon final states from 1.4×10^{11} positive and 2.9×10^{10} negative 209-Gev muons in a magnetized iron calorimeter has set a lower limit of 9 Gev/c^2 on the mass of a heavy neutral muon (M^0), and a 90%-confidence level upper limit of $\sigma(\mu N \to b\bar{b}X)B(b\bar{b} \to \mu X) < 2.9 \times 10^{-36}$ cm^2 for the production of bottom hadrons by muons. The dimuon mass spectrum from 102,678 trimuon final states places a 90%-confidence level upper limit for the muoproduction of upsilon states: $\sigma(\mu N \to \mu \Upsilon X)B(\Upsilon \to \mu^+\mu^-) < 22 \times 10^{-39}$ cm^2 . In addition, analysis of 71 rare multimuon events, including 4- and 5-muon final states, is presented.

I. INTRODUCTION

Much of particle physics appears to be described by gauge theories. The standard model[1] is based on the group $SU(3) \times SU(2) \times U(1)$, spontaneously broken into $SU(3)_c \times U(1)_{em}$. This theory was elaborated by the work of Glashow, Illiopoulos and Maiani[2], which explained charmed hadrons. This, in turn, was naturally extended by Kobayashi and Maskawa[3] to 3 left-handed doublets of quarks, which allowed the incorporation of the τ lepton and its neutrino, and the new bottom quark which comprises the Υ family[4]. If this model is to form the bulwark of our understanding of the structure of matter, then it must be comprehensively studied.

This exploration may proceed down several avenues. One can look for currents which have not been seen, but which have not been experimentally ruled out. A current of this type is a right-handed weak current coupling the muon to a neutral heavy muon. Another route is to study the interactions of the newly discovered quark to see if it behaves in a manner analogous to the lighter and better studied quarks. The experimental study of hadrons with bottom quarks is just beginning. The primary experimental evidence involves the detection of the direct leptons from semileptonic decays of bottom mesons[5]. A third approach is to look for rare or "exotic" phenomena. A rich source of such phenomena is multimuon final states. There have been reports of "super" neutrino-induced trimuon events at Fermilab[6], which are not consistent with the conventional physics usually employed to explain these trimuons. In addition, experiments at CERN[7] and Fermilab[8] have observed neutrino induced 4-lepton events for which an adequate explanation is lacking.

A particularly fertile ground for the exploration of these areas is

muon physics. The right-handed chirality of a high energy muon beam provides a unique probe of the right-handed weak current. As a source of virtual photons, the muon beam can explore the behavior of heavy quark states in kinematic regions inaccessible through other means. Finally, by taking advantage of the ability of muons to penetrate vast quantities of matter, one can use massive targets to conduct searches for rare processes with cross sections as low as $10^{-39}cm^2$.

For these purposes, a Fermilab muon experiment, E203/391, was performed to study a broad range of muon-induced physics. The Berkeley-Fermilab-Princeton multimuon spectrometer was designed to have a high sensitivity to any number of muons in the final state. A large solid iron magnet integral with the target provided uniform acceptance over the entire length of the apparatus. The experiment was unique in its ability to do multimuon physics because of its full acceptance over its entire fiducial region, due to the lack of any insensitive area in the vicinity of the muon beam.

This thesis presents results from data taken with the multimuon spectrometer in the first half of 1978. Chapters II and III describe the experiment and its analysis. Chapter IV presents a search for heavy neutral muons. Chapters V and VI detail limits on the muoproduction and virtual photoproduction of bound and open bottom quark states. Chapter VII shows the analysis of the sample of 71 rare multimuon final states. Descriptions of the results in chapters IV, V, and VI have appeared in references 9, 10 and 11, respectively.

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II. THE EXPERIMENT

A. The Muon Beam

The muon beam was produced by the decay in flight of pions and kaons produced by the 400 Gev proton beam incident on a 30 cm aluminum target. Figure II.1 shows a schematic diagram of the Fermilab muon beam. A series of quadrupole magnets, labelled Q1, focussed the secondaries from the target into a 400 m long decay pipe. Momentum selection was accomplished by bending the beam to the right with dipole D1 and then to the left with dipole D2. The currents in these dipoles were set to select a particle of one sign and a momentum near 215 Gev/c. The momentum acceptance was 2.5%. The 60 feet of polyethylene absorber in dipole D3 stopped hadrons in the beam. Quadrupole Q4 focussed the beam on the apparatus, while dipole D4 bent the beam into the Chicago cyclotron magnet (CCM) for targetting on the spectrometer.

Figure II.2 shows the beam line and its monitoring from the focussing quadrupoles to the multimuon spectrometer. Hodoscopes and proportional wire chambers before and after the dipole magnets and the Chicago cyclotron magnet identified beam particles and provided momentum measurements. Multiple coulomb scattering of muons in the polyethylene and muons scraping the beam elements produced halo muons in the muon laboratory. Several veto counters and a large veto wall identified these halo muons. The number of muons in the halo was roughly equivalent to the number of muons in the beam. The muon beam produced intensities up to a total 6×10^6 muons/spill in the beam area, which was 8 inches high by 13.5 inches wide at the front of the spectrometer. The yield of total beam muons per proton was as high as 4×10^{-7}.

B. Multimuon Spectrometer

A schematic view of the multimuon spectrometer is shown in figure II.3. It is composed of 91 plates of steel 4 inches thick and 8 feet square. Each plate has 2 slots cut in it through which 2 coils running the length of the spectrometer were placed. The fiducial area, located between the coil slots, was magnetized to a total 19.7 kG vertical field, which was uniform to 3% over the central 1.4x1 m area of each slab.

The steel slabs were distributed with one lone plate in front followed by groupings of five slabs, called modules. An individual module is shown in figure II.4. Modules were separated from each other by a 10 inch gap. The first slab and the slabs in the first 15 modules served as the target with a density of 6.1 kg/cm^2. The steel also served as a hadron and photon filter with an average density in the spectrometer of 4.7 gm/cm^3. Particles were required to traverse 4 modules, almost 12 absorption lengths, before identification as muons.

Three types of magnetic measurements were made to determine the magnetic field in the multimuon spectrometer. Flux loop measurements determined the absolute normalization for the field integrals in the various modules. These were done with wire loops around the steel plates that measured the induced EMF as the magnet was ramped on and off. Search coil measurements in the gaps between iron slabs determined the relative field shape as a function of x and y. Finally, various physical measurements necessary to calculate the field integral were performed, such as determining the width of iron in each module. The field was mapped with 0.2% accuracy in the central area of the spectrometer. The

polarity of the field was reversed periodically.

Hadron showers produced in interactions were sampled every 10 cm by plastic calorimeter scintillation counters placed after every slab in the first 15 modules. The calibration of the calorimeter was obtained by statistical comparison with the magnetic measurement of the energy lost in an interaction (subtracting the outgoing muon energies from the energy of the incoming muon). The rms accuracy of the hadron calorimetry was $\Delta E = 1.5E^{\frac{1}{2}}$ for ΔE and E in Gev, with a minimum uncertainty of 2.5 Gev.

After every even-numbered module, beginning with the fourth, banks of scintillation trigger counters were installed. The configuration of these counters is shown in figure II.5. They consist of 4 large paddle counters at the top and bottom, and six narrow staves in the middle, framed by two wider staves.

C. Wire Chambers

A multiwire proportional chamber was placed after every module and the single slab at the front. The proportional chambers had three planes of wires. There were 336 anode wires spaced at 3 mm which read out coordinates in the horizontal (x), or bend plane, direction. Coordinates in the diagonal (u) and vertical (y) directions were registered by by means of 5 mm wide cathode strips composed of 4 high voltage wires apiece. The diagonal plane consisted of 176 such strips and the vertical 192. Each strip was connected to one input of a differential amplifier as shown in figure II.6. Although spread over many cathode strips, the induced charge produced a count only in the one or two electronics channels

closest to the peak, even when the pulse height far exceeded threshold. This center-finding circuitry gave twice as good a resolution as that achievable with conventional circuitry. The separation between the diagonal and vertical cathode planes and the anode plane was 1 cm. The chambers were active over the entire fiducial area 1.8 m high by 1.1 m wide.

The resolution of the anode plane (x) measurements was 1 mm and the resolution of the cathode plane (u and y) measurements was 3 mm. Outside the beam region the anode and cathode planes had efficiencies of 95% and 94% respectively. In the central beam region at the highest beam intensities, these efficiencies for the most upstream chambers could drop as low as 83% and 59%. Generally, chambers would have efficiencies down to 88% for the anode plane and 76% for the cathode planes in the central beam region at highest beam flux. Data from the chambers was read out for 70 nsec during a trigger.

Attached to every multiwire proportional chamber was a single drift chamber plane with 56 vertical wires measuring coordinates in the bend plane. The drift cell width was 3/4 inch and the distance from the sense wires to the field-shaping high voltage plane was 1/8 inch. Each drift chamber covered the entire fiducial area. The drift chambers were gated for 250 nsecs during a trigger. The resolution of each drift chamber was 250 microns and their average efficiency was 98%. The drift chambers provided the maximum resolution compatible with multiple coloumb scattering in the bend plane in order to produce more precise muon momentum determination. The drift chamber system is described in detail in Ref. 1.

D. Triggers

The apparatus ran with four simultaneous triggers: "beam", "one muon", "two muon", and "three muon". The "beam" trigger required a muon to trigger in the beam hodoscope counters upstream of the spectrometer without any of the halo veto counters firing. This trigger was always used in coincidence with all other triggers and provided a trigger by itself when prescaled by 3×10^5. The "one muon" trigger was used to detect high Q^2 muon scattering and therefore required each of three consecutive trigger banks to have a hit in a paddle counter and to have no hits in any stave.

The "two muon" trigger required 3 trigger banks to have ≥ 2 hits and at least 20 Gev of energy deposited in the calorimeter. In addition, the hits in the most downstream contributing trigger bank were required to be non-adjacent. This trigger is described in detail in Ref. 2. The "three muon" trigger required three consecutive trigger banks to have ≥ 3 hits, but did not involve the calorimeter. It also demanded that one of the hits be non-adjacent to the other two hits in the most downstream two trigger banks. The rates of the "one", "two" and "three muon" triggers relative to one beam muon were 3×10^{-6}, 8×10^{-6}, and 1.2×10^{-5}, respectively.

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Chapter II

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Figure Captions

Figure II.1. Schematic diagram of the Fermilab Muon beam from the extracted proton beam through the Chicago cyclotron magnet (CCM) just upstream of the multimuon spectrometer.

Figure II.2. Detailed view of the beam magnets, proportional chambers and scintillation counters along the muon beam in enclosures 103 and 104 and in the muon laboratory.

Figure II.3. Schematic view of the apparatus. S_1-S_{12} are trigger scintillators (1 of 8 banks). DC and PC are 1 of 19 pairs of drift and proportional chambers. Each proportional chamber measures projections on three coordinates. The scintillators labelled SC are 5 of 75 counters performing hadron shower calorimetry.

Figure II.4. Side view of one module containing 5 steel plates followed by 5 calorimeter counters and the trigger scintillator bank, proportional chamber and drift chamber in the large gap that separates the groups of 5 plates.

Figure II.5. An exploded view of the detectors within a typical gap between magnet modules. The trigger hodoscope follows the calorimeter counter. Counters S_1, S_2, S_{11} and S_{12} are "paddles" 20.75 inches wide and 23.8 inches high. Counters S_3-S_{10} are "staves". S_3 and S_{10} are 41.5 inches wide and 5.98 inches high while S_4-S_9 are 41.5 inches wide and 1.55 inches high.

Figure II.6. The Network of differential amplifiers sensing the center of the charge distribution induced on the proportional chamber cathode strips.

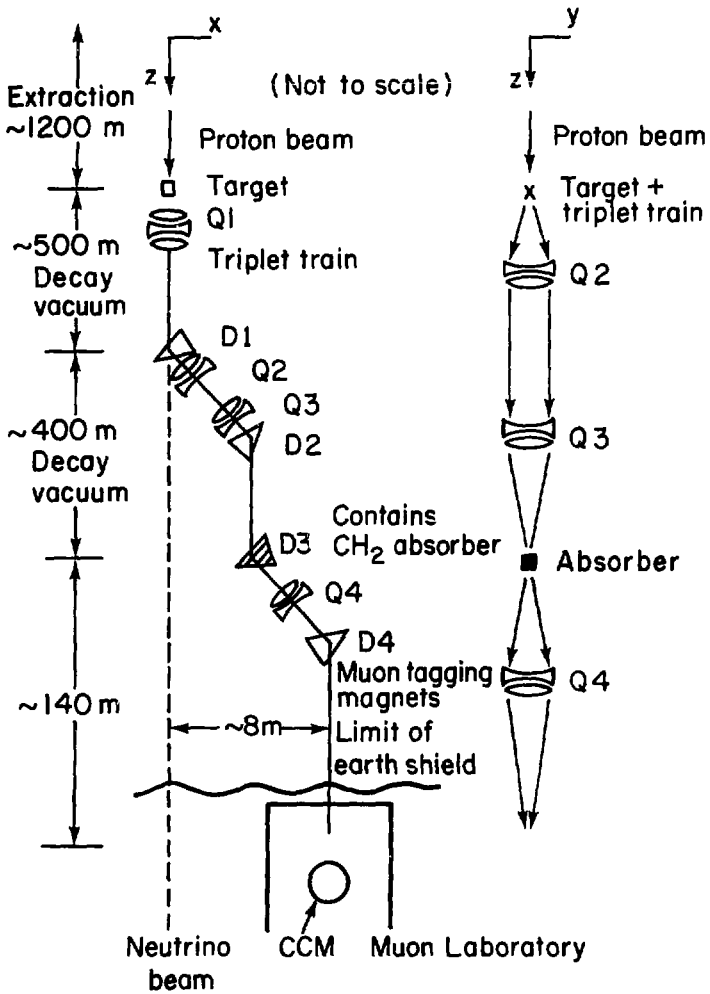

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Figure II.1

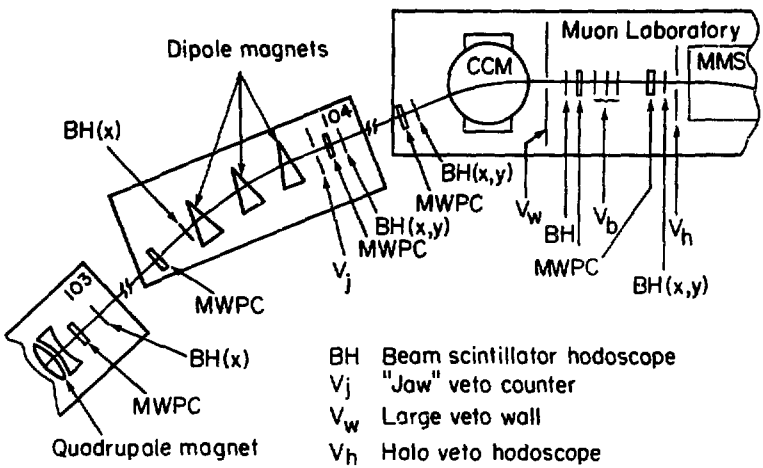

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Figure II.2

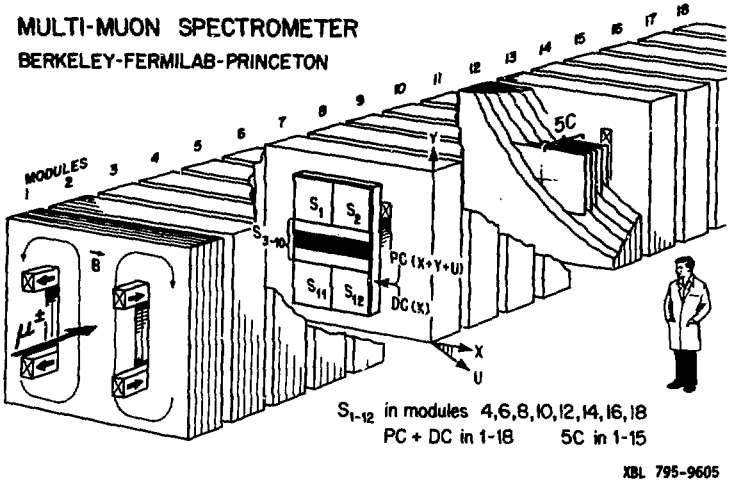

MULTI-MUON SPECTROMETER
BERKELEY-FERMILAB-PRINCETON
MODULES
1
2
3
4
5
6
7
8
9
10
11
12
13
14
15
16
17
18
5C
Y
X
U
S_1
S_2
S_{3-10}
S_{11}
S_{12}
PC (X+Y+U)
DC (X)
$\vec{B}$
$\mu^{\pm}$
S_{1-12} in modules 4,6,8,10,12,14,16,18
PC + DC in 1-18 5C in 1-15
XBL 795-9605

Figure II.3

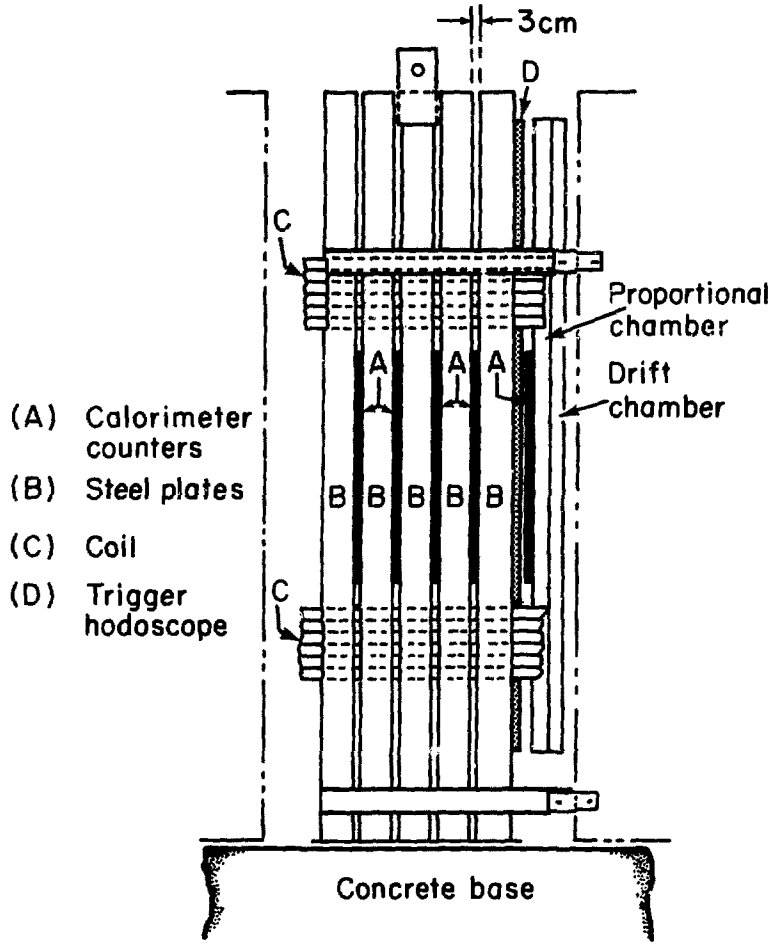

3cm
D
C
Proportional chamber
A
A
A
Drift chamber
(A) Calorimeter counters
(B) Steel plates
(C) Coil
(D) Trigger hodoscope
B
B
B
B
B
C
Concrete base
XBL8010-2254

Figure II.5

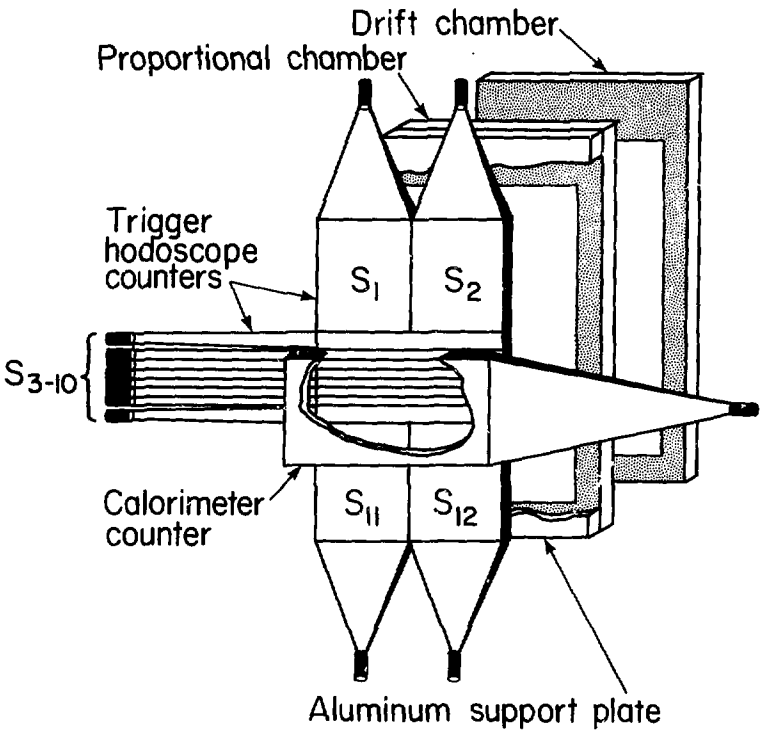
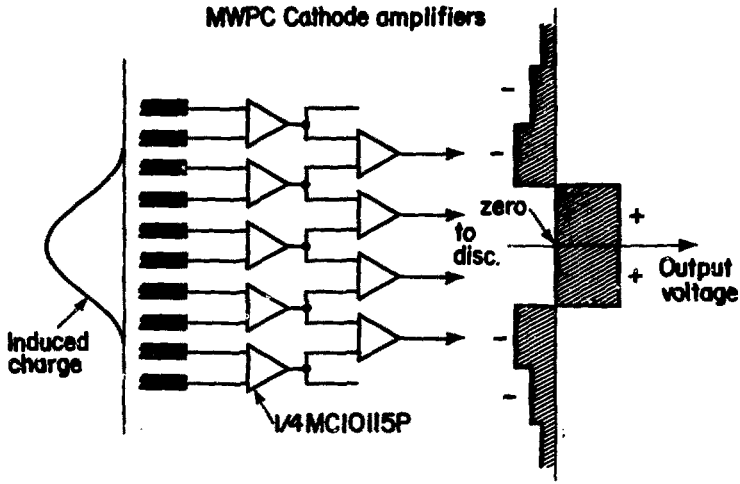

Drift chamber
Proportional chamber
Trigger
hodoscope
counters
S1
S2
S3-10
Calorimeter
counter
S11
S12
Aluminum support plate
XBL 795-1602

Figure II.5



XBL 795-1599

Figure II.6

III. ANALYSIS

A. Track Finding

The track finding program combines contiguous proportional chamber hits into single hits with measurement errors equal to $1/\sqrt{12}$ the distance between the first unstruck wires on either side of the group of wire hits. If a diagonal (u) plane wire is struck within .75 cm of a hit x-wire and hit y-wire crossing, the x, y, and u hits are declared a matched triplet. The program begins at the back of the spectrometer and requires three triplets or two triplets and unmatched x and y hits in a third chamber. The three chambers containing these hits must not be separated from each other by more than one empty chamber.

The track is extended one chamber at a time, where a new triplet or unmatched hits are attached, the trajectory is recalculated, a projection of the track is extended into the next chamber and a window for searching for new hits is opened in this chamber. This procedure continues until the track finder passes two contiguous chambers where the search window contained no hits or the location along the beam (z) axis of the event vertex determined by calorimetery is reached.

B. Calorimeter Vertex

There are two methods of searching for the location along the z axis for the event vertex by examining the pulse heights in the calorimeter counters. In the case of a "one muon", or "two muon" trigger, or a "three muon" trigger accompanied by more than 40 GeV of energy deposited in the calorimeter, an "inelastic" calorimeter vertex is found. In the other cases, an "elastic" calorimeter vertex is found. If the inelastic

vertex finder fails on a "three muon" trigger, the elastic vertex finder is used. In all other cases, if the vertex finder fails the vertex is set at the front of the spectrometer so as to not interfere with track-finding.

The elastic calorimeter vertex finder computes the likelihood of the vertex in each steel plate using normalized 1 and 3 particle calorimeter distributions. The routine uses the pulse heights from all the calorimeter scintillators in the calculation and searches from the first plate to the plate before the most downstream trigger-scintillator bank contributing to the event trigger. The inelastic calorimeter vertex finder searches for the calorimeter counter with the largest pulse height. It then computes for each slab the difference between the number of upstream counters with less than and with greater than 8% of this pulse height. The vertex is assigned to the slab with the maximum value of this difference.

C. Beam Track Finding

The information from the wire chambers, shown in figure II.2 along the muon beam lines in enclosures 103, 104 and the area upstream of the multimuon spectrometer in the muon laboratory is used with the first proportional chamber in the spectrometer to determine the slope, position, momentum and their errors for the incident beam muon at this first chamber. The momentum is measured from the bend of the dipoles in enclosure 104 and the Chicago cylotron magnet in the upstream end of the muon laboratory. If the chi-square for this fit is poor, the chamber contributing the largest residual is discarded and the track is refit.

Irrespective of its χ^2, the fit muon trajectory is then projected into the spectrometer, one chamber at a time, and triplets or, if there are none, unmatched hits are assigned to the track. The trajectory is then refit using the new chamber hits and projected into the next chamber. The procedure continues until the calorimeter vertex is reached, or in the case of a failed calorimeter vertex in the first slab, until the most downstream trigger bank contributing to the event trigger.

After all track finding is complete, the two drift chamber hits closest to the fit proportional chamber trajectory in the x view are attached to every track. The choice of which of these hits, if any, to incorporate in the track is made by the track fitting routine.

D. Track Fitting

The track fitting program begins with the track provided by the track finding program. At first, only proportional chamber tracks are fit. Once a track has been fit in the bending plane, the program scans the drift chamber track arrays and replaces proportional chamber hits with chosen drift chamber hits if they lie within a distance equal to three times the uncertainty in the position of the fit track. The combined drift and proportional chamber hits are then fit by the momentum fitting routine again.

E. Momentum Fitting Routine

For outgoing tracks, the momentum fitting routine takes as input a point along the z axis for reference and all the proportional and drift chamber hits downstream of that point. It makes a simultaneous fit to

the free parameters describing the muon tracks. In the bending plane, these are the transverse position x_0 and direction tangent s_0 of the muon at the reference point and the muon momentum $p = 1/\rho_0$, projected in that plane.

N additional free parameters d_j are introduced equal to the projected transverse momentum impulse due to multiple coulomb scattering in each of the N magnet segments that the muon traverses after the reference point. A magnet segment is defined as the steel between the n wire chamber hits located at x_i. Thus, there are N additional measurements d_j with variances e_j, where e_j is the rms value of d_j appropriate to the thickness of the iron segment. When the d_j are introduced, σ_i, the errors on the x_i, become deviations due only to intrinsic chamber measurment error.

Each magnet segment imparts an impulse h_j of transverse momentum to the muon. The h_j were corrected for departure from normal incidence. In addition, the measured coordinate X_i was given a correction ΔX_i for the effect of muon energy loss in each magnet segment. Each iteration of the fit changed these ΔX_i appropriately, based on the last best fit momentum. Hence the full chi-squared is

$$\chi^2 = \sum_{i=1}^{n} \frac{(x_i - (X_i + \Delta X_i))^2}{\sigma_i^2} + \sum_{j=1}^{N} \frac{d_j^2}{e_j^2}$$

where

$$x_i = x_0 + s_0 z_i + \sum_{\substack{j=1 \\ w_j<z_i}}^{N} (z_i - w_j)(h_j\rho_0 - d_j\rho_j)$$

$$\Delta X_i = \sum_{\substack{j=1 \\ w_j<z_i}}^{N} (z_i - w_j) h_j \Delta\rho_j$$

and where z_i and w_j are the coordinates along the beam axis relative to the reference point of the measurement planes and magnet segment midpoints, and $\rho_j=\rho_o+\Delta\rho_j$, where $\Delta\rho_j$ is produced by the energy loss in the iron.

The best fit to the free parameters x_o, s_o, ρ_o, and $(d_1,\ldots,d_N)$ was obtained by solving the N+3 simultaneous linear equations

$$\frac{\partial\chi^2}{\partial x_o} = \frac{\partial\chi^2}{\partial s_o} = \frac{\partial\chi^2}{\partial \rho_o} = \frac{\partial\chi^2}{\partial d_1} = \ldots = \frac{\partial\chi^2}{\partial d_N} = 0.$$

For the non-bending plane fit (y coordinate) to an outgoing track the momentum is taken from the x fit and is not a free parameter. For beam or incoming tracks fit in the spectrometer, the incident angle and direction in the x and y views is taken from a fit made to the beam system.

If the momentum is being fit as a free parameter, then the routine iterates using as input to the fit a value of the momentum that is a function of the previous guessed input values and output values returned by the routine. For all tracks and views, if the chi-square of the fit track is unsatisfactory the routine removes the measurement plane whose hit contributes the largest amount to the chi-square and refits the tracks. No more than 1/3 of a track's hits may be removed and a minimum of 5 hits must remain. In the bending view each measurement plane may contain 2 drift chamber and one proportional chamber hit for each track. The fitting routine tries swapping the chosen hit for another before it removes the measurement plane. The fit momentum resolution is 8%.

F. Vertex Finding

In preparation for vertex finding, the routine eliminates tracks

that penetrated but were not detected by the trigger counters. These tracks are due to muons out of time with the event by more than the 19 ns r.f. bucket time. Tracks that were mistakenly broken into two or more segments by the track-finder program are rejoined by composing a new track from hits in the segments. Tracks are selected for rejoining on the basis of the number of hits they have in common and the angle they make with each other at their point of contact or closest approach. The decision to merge tracks is based on the χ^2 of a fit made to a track composed of the combined hits of both tracks. Single tracks that the track finder reconstructed as two tracks have one of the duplicates removed.

Finally, tracks with over 5 blank measurement planes between their apparent termination and their fit exit from the spectrometer in either the x or y view are eliminated. The event is thrown out if no secondary tracks remain, or, in the case of a "two muon" or "three muon" trigger, if less than two secondary tracks remain. These are events which accidently triggered as having two or more secondary tracks when these tracks did not actually occur in the event.

The vertex finder first chooses the secondary tracks to be used in determining the vertex on the basis of their distance of closest approach to the beam track, the error in this distance, the chi-square of their original fit and the distance they extend upstream of their point of closest approach. The z position of the vertex, zv and its error σ_{zv}, are then chosen by a weighted average of the included tracks' closest points of approach and the calorimeter vertex if the chi-square per degree of freedom of the fit including it with the track vertex is

less than 3.

The vertex finder scans $5\sigma_{zv}$ on either side of zv in 10 cm steps, using the fits of the included tracks, the beam track, and their errors to determine the most likely point (x_0, y_0) in common for all these tracks at each step. A chi-square is determined for each point, where, given a step in z, the index runs over the included tracks:

$$\chi^2 = \sum_i \frac{(x_i - x_0)^2}{(\Delta x_i^2 + \Delta x_0^2)} + \frac{(y_i - y_0)^2}{(\Delta y_i^2 + \Delta y_0^2)} .$$

The minimum chi-square determines the z position of the vertex. The vertex finder then performs a 1 cm scan in a 20 cm range centered on this vertex, finds a new best vertex and finally performs a 1/3 cm scan in a 2 cm range centered on this vertex.

During vertex finding procedures the calorimeter vertex is examined for consistency with the track vertex. The calorimeter vertex is considered consistent if it is within a distance, equal to 1.5 times the uncertainty in its position, away from the vertex determined by the tracks and calorimeter vertex combined. If it is found consistent, it is included with its error in the chi-square scan. If it is not, it is removed and the vertex finding begins again without it. If the inelastic calorimeter vertex is available, then the vertex finder does a 1 cm scan in a 100 cm range centered on the calorimeter vertex and is not allowed to discard the calorimeter vertex. The 1/3 cm scan follows as before. If the overall chi-square for the vertex is unsatisfactory, the routine attempts to throw one or more tracks out of the set of included tracks and repeats the entire procedure.

Once this vertex has been determined, it is attached to all tracks

and they are then refit by the momentum fitting routine. If any track except the beam track has a large chi-square from this fit, its original fit is restored and it is considered excluded from the vertex. The severity of the chi-square cut is adjusted to provide a sample of at least 3 outgoing tracks or 2 outgoing tracks and an inelastic calorimeter vertex to be attached. However, a track is never included in the vertex if its chi-square per degree of freedom exceeds 7.5 in either x or y view when the vertex is attached.

If it is found that the sample of tracks attached to the vertex is not the same as that used in previously determining the vertex or that any measurement planes were removed in the momentum fit with the vertex attached that were included in the original momentum fit, the tracks are all refit without the vertex attached, but with all the newly removed measurement planes on each track removed a priori. The entire vertex scanning and determining procedure is then repeated. If is found that the use of an inelastic calorimeter vertex resulted in too large a chi-square, the vertex finding and fitting procedure is repeated with the calorimeter vertex treated as though it were an elastic vertex. Once the new vertex has been determined, all these tracks are once again fit with this vertex included as one of their hits and they are constrained to go through it.

G. Acceptance Modeling

Monte Carlo calculations of the detector acceptance are based on a standard program onto which the various physics generators are coupled. These generators include the muoproduction of neutral heavy muons, psis, upsilons, pions, kaons, charmed mesons, and bottom mesons. The Monte

Carlo program uses a sample of real beam muon events to simulate the real beam distribution. These beam muons are propagated through the spectrometer to the interaction vertex.

The daughter muons from the generator are propagated until they leave the spectrometer. This propagation includes energy loss from μ-e collisions, muon bremsstrahlung and electron pair production. It also calculates the bending of muon trajectories in the magnetic field and includes multiple coulomb scattering. Large angle scattering is parameterized by a nuclear form factor. A basic attempt is also made to model the hadronic shower spread through the chambers. The Monte Carlo also produces calorimeter pulse heights and trigger counter latches. Interactions that trigger the apparatus are written on tape using the same format employed in actual data taking.

IV. LOWER LIMIT ON NEUTRAL-HEAVY MUON MASS

A. Experimental and Theoretical Background

Considerable speculation has been devoted to the possible existence of heavy neutral gauge leptons. Variations of the standard SU(2)xU(1) model[1] have been proposed which include[2] M^0's. Grand unification schemes frequently introduce M^0's, e.g. those[3] which embed $SU(2)_L xU(1)_R$ in $SU(3)_L xSU(3)_R$. In addition to the M^0, heavy doubly charged gauge muons (M^{++}) have been proposed in the context of an extended SU(2)xU(1) theory in doublets with the known singly charged leptons[2].

There exist few experimental limits on the masses of heavy muons. Studies of π and K decay[4] exclude the M^0 mass from the range $m_\mu < m_{M^0} < m_K$. A bubble chamber study of ν_μ-N interactions[5] sets a 90%-confidence lower limit of 1.8 GeV/c^2 on the mass of the heavy muon M^-. Although there are 90%-confidence lower limits of 2.4 GeV/c^2 from ν_e-N scattering[6] and 8.4 GeV/c^2 from ν_μ-Fe interactions[7] on the M^+ mass, there is no further experimental constraint on the M^0 mass.

Possible evidence for M^0 production has arisen from three experiments. Two μ^-e^+ events produced by ν_μ-N interactions below 30 GeV in the SKAT bubble chamber[8] were attributed[9] to the production of an M^0 with $1.4 < m_{M^0} < 2.4$ GeV/c^2. However, no corroborating evidence for the M^0 has resulted from the study[10] of ν and $\bar{\nu}$ induced μe pairs. In a cosmic ray experiment[11] deep underground, five events were interpreted either as evidence for a heavy lepton with mass 2-4 GeV/c^2 or as the cascade[12] of a new charged heavy lepton to an M^0. However, two subsequent searches[13] found no such events. Originally the observation of neutrino-induced

trimuon events at Fermilab[14] prompted their interpretation[15] as examples of M^0 production. Further experiments and analyses found this phenomenon to be compatible with conventional processes: heavy lepton production could account for no more than 10-20% of these events[16].

B. Rate Calculation

We have calculated the expected rates for $\bar{M}^0$ and M^{++} production in this experiment, assuming the incident muon to be coupled with Fermi strength to the M by means of a right-handed weak current. The right-handed coupling, present in most models containing a heavy gauge lepton, is compatible with our experimental conditions due to the $\sim$80% left-handed polarization of the μ^+ beam[17]. In the limit of negligible muon mass, invariance to weak isospin rotation gives $\sigma(\mu^-(L.H.)N\rightarrow\nu_\mu X)=\sigma(\nu_\mu N\rightarrow\mu^- X)$, where L.H. refers to the left-handed muon helicity and N is an average of proton and neutron. Also, for negligible M^0 mass, $\sigma(\mu^-(L.H.)N\rightarrow M^0X)=(g_L/g)^2\sigma(\mu^-(L.H.)N\rightarrow\nu_\mu X)$, where g_L^2/g^2 is the ratio of left-handed coupling strengths for M^0 and ν_μ. Finally, $\sigma(\mu^+(L.H.)N\rightarrow\bar{M}^0X)=(g_R/g_L)^2\sigma(\mu^-(L.H.)N\rightarrow M^0X)$, where g_R^2/g_L^2 is the ratio of abnormal-helicity to normal-helicity weak coupling strengths[18] for the M^0. For a right-handed current of Fermi strength this ratio is unity. Except for effects of finite lepton mass, these equations combine to give $\sigma(\mu^+(L.H.)N\rightarrow\bar{M}^0X)=(g_R/g)^2\sigma(\nu_\mu N\rightarrow\mu^- X)$.

Using the simplest parton model with single W^+ exchange[19], invoking the Callan-Gross relation[20] and considering only $\Delta S=\Delta C=0$ processes and isoscalar targets,

$$\frac{d^2\sigma(\mu^+(L.H.)N\rightarrow\bar{M}^0X)}{dvdy} = \left(\frac{g_R}{g}\right)^2 \frac{G^2 E m_N F_2(x)}{\pi y}$$

where $v=xy=Q^2/s$, $(1-y)$ is the fraction of the laboratory muon energy re-

tained by the $\bar{M}^0$, and $F_2(x)=18\nu W_2^{\gamma N}(x)/5$. We parameterize $\nu W_2^{\gamma p}$ as in Ref. 21 and set[22] $\nu W_2^{\gamma N}=(1-0.4x)\nu W_2^{\gamma p}$. The differential cross section is independent of $\bar{M}^0$ mass, except for kinematic restriction of the allowed area of the Q^2-ν plane.

The differential decay rate for $\bar{M}^0\to\mu^+\mu^-\bar{\nu}_\mu$, where the $\bar{M}^0$ is coupled to the μ^+ by a (V+A) current, is

$$\frac{d^5\Gamma(\bar{M}^0\to\mu^+\mu^-\bar{\nu}_\mu)}{dx_-dx_\nu d\phi_\nu d\cos\Theta_\nu d\phi_-} \propto x_\nu(1-x_\nu)(1-h\cos\Theta_\nu)$$

In the $\bar{M}^0$ rest frame $x_-(x_\nu)$ is $2p/m_{\bar{M}^0}$ for the $\mu^-(\bar{\nu}_\mu)$, Θ_ν and Θ_ν define the $\bar{\nu}_\mu$ direction relative to the $\bar{M}^0$ direction, Θ_- and ϕ_- define the μ^- direction relative to the $\bar{\nu}_\mu$ direction, and h is the $\bar{M}^0$ helicity. Since the $\bar{M}^0$ carries the left-handed polarization of the incident μ^+, the two muons are emitted preferentially forward and together carry an average of 80% of the $\bar{M}^0$ energy in the laboratory.

C. Results

Monte Carlo events have been generated according to the above formulae at lepton masses of 1,2,3,5,9,12 and 14 GeV/c^2. Simulated $\bar{M}^0$ and M^{++} events at each mass are binned in $\sqrt{Q^2}$ and in $p_\perp$, the daughter muon momentum transverse to $\vec{Q}$. For this analysis, Q^2 is defined by taking the highest-energy beam-sign final state muon to be a scattered beam muon. The $\bar{M}^0$ (M^{++}) Monte Carlo events are compared to data events containing exactly two opposite- (same-) sign reconstructed final-state muons. The data events consist of 76,350 opposite-sign and 46,615 same-sign dimuon final states produced by 1.4×10^{11} positive and 2.9×10^{10} negative 209-GeV muons.

Kinematic cuts were chosen individually for each heavy lepton type and mass in order to exclude data while retaining Monte Carlo $\bar{M}^0$ or M^{++} events. Primarily, these cuts demand a particular range of invariant mass[23]. In addition, for m_{M^0} >3, >2, or <3 GeV/c^2, respectively, the cuts require a 9 GeV minimum outgoing muon energy, a -5 GeV minimum missing energy, or a 50 GeV minimum ν. The cuts suppress the principal backgrounds of charm production and π- and K-decay. An empirical contour then was drawn for each $\sqrt{Q^2}$-$p_\perp$ plot in order to contain all the data events on the low $p_\perp$, low $\sqrt{Q^2}$ side. The same contour was drawn on the corresponding plot for simulated M events. (If the same contour[24] and cuts, except for the dimuon mass cut, were used for all masses, the limits presented below would rise by a factor of 1.6 on the average). Figure IV.1 shows the plots and contour for data and Monte Carlo corresponding to 6 GeV/c^2 $\bar{M}^0$ production. The Monte Carlo event populations on the high $p_\perp$, high $\sqrt{Q^2}$ side of the contours then provide the cross section limits.

Figure IV.2 displays the mass-dependent limits on the product of cross section and $\mu\mu\nu$ branching ratio (σB) for $\bar{M}^0$ and M^{++} production. Also indicated are the calculated σB for the production of $\bar{M}^0$'s and M^{++}'s, where the branching ratio is assumed to be 0.1 and 0.2 for $\bar{M}^0$ and M^{++}, respectively. At 90% confidence the data exclude the production of a $\bar{M}^0$ or M^{++} coupled with Fermi strength to a right-handed current in the mass range $1<m_{M^0}<9$ GeV/c^2. Without a special mechanism to suppress pair production, doubly-charged leptons in this mass range would have been detected at PETRA. No comparable limits on $\bar{M}^0$ production in this range are available from any other experiment.

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Chapter IV

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24. The contour outlines the region $\sqrt{Q^2} > 3.0$ and $p_\perp > 3.44 - 0.11\sqrt{Q^2}$.

Figure Captions

Figure IV.1. Two-dimensional event distributions vs. $\sqrt{Q^2}$ and $p_\perp$, defined in the text. The vertical scale is logarithmic; bin populations range from 0 to 450. Distribution (a) shows the data and the empirically chosen contour within which these events are contained. Distribution (b) is 77.4× the simulated population from production and decay of a 6 GeV/c^2 $\bar{M}^0$, with the assumptions described in the text. The events in (b) lying outside the contour in (a) give the quoted σB limit at this mass.

Figure IV.2. Experimental upper limits and calculated cross section-branching ratio products σB for heavy-muon ($\bar{M}^0$ and M^{++}) production by 209-GeV muons, plotted vs. heavy muon mass. The calculation assumes $B(M \to \mu\mu\nu)=0.1$ ($\bar{M}^0$) or 0.2 (M^{++}), and right-handed coupling of μ^+ to M with Fermi strength ($g_L=g_R$).

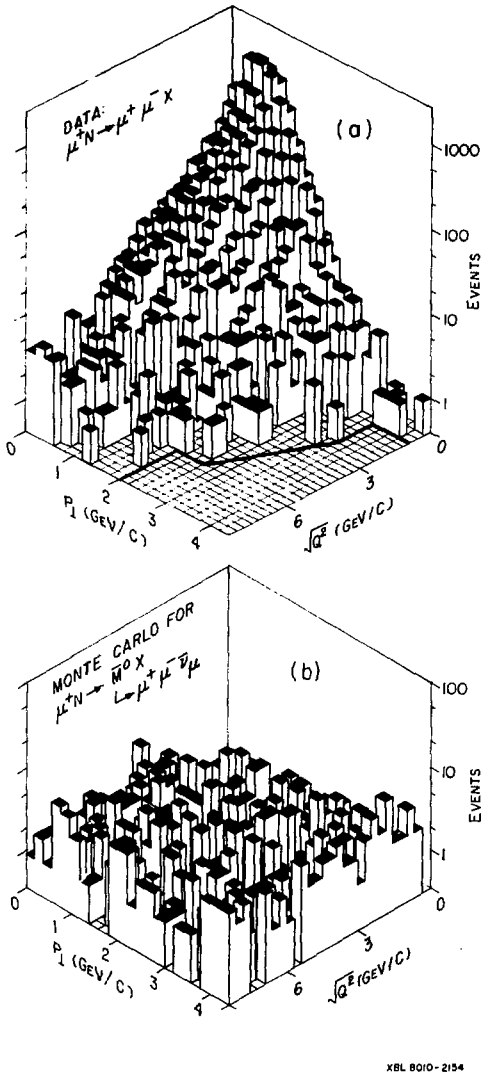
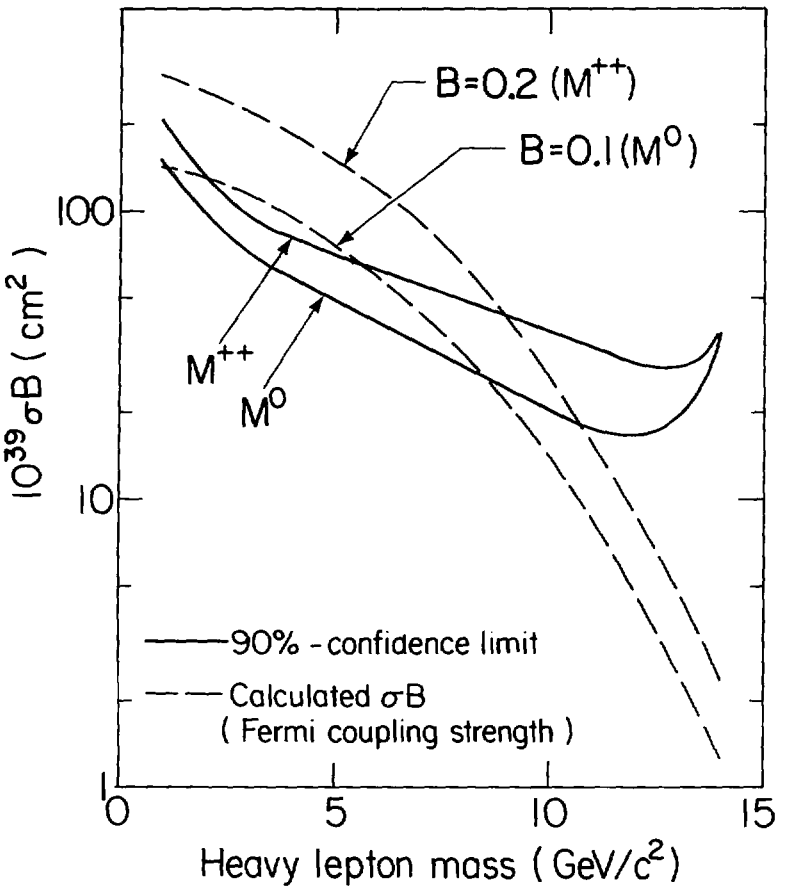

DATA:
$\mu^+N \to \mu^+\mu^- X$
(a)
1000
100
10
1
EVENTS
0
1
2
3
4
$P_\perp$ (GEV/C)
0
3
6
$\sqrt{Q^2}$ (GEV/C)
MONTE CARLO FOR
$\mu^+N \to M^0 X$
$\hookrightarrow \mu^+\mu^-\bar{\nu}_\mu$
(b)
100
10
1
EVENTS
0
1
2
3
4
$P_\perp$ (GEV/C)
0
3
6
$\sqrt{Q^2}$ (GEV/C)
XBL 8010-2154

Figure IV.1



XBL 809-2016

Figure IV.2

V. A LIMIT ON T MUOPRODUCTION

The dimuon mass spectrum from an integrated luminosity of $0.78 \times 10^{39} cm^{-2}$ is derived from 102 678 trimuon final states. This data sample contains 6693±355 examples of J/ψ and ψ' production[1] and contains invariant masses up to 11.5 Gev/c^2. In every event, all three outgoing muons are fully momentum-analyzed and are subjected to an energy-conserving one-constraint fit using calorimetric measurement of the associated shower energy. The quality, statistical power and range of this sample make it exceptionally suitable for an investigation of the virtual photoproduction of heavy quark states by muons. At present, there is no other comparable sample from any other experiment. We have chosen here to use the sample to search for muon-induced virtual photoproduction of T states.

No limit on T production by real or virtual photons has been published. A conference report[2] based on results from the Bologna-CERN-Dubna-Munich-Saclay (BCDMS) experiment presents the limit $\sigma(\mu N \to TX)B(T \to \mu^+\mu^-) < (6\pm3) \times 10^{-39}$ cm^2 (90% confidence) for ~275-GeV muons, where the error is systematic. This limit is based on 761 multimuon events corresponding to an integrated luminosity[2] of 0.7×10^{39} cm^2 . A third muon was observed in 11% of these events. No calorimetric information was available. With 48% T acceptance, the BCDMS limit corresponds to ≤2 T candidates (90% confidence). In total, the experiment observed 24 events between 8 and 12 GeV/c^2 in dimuon mass. These were compared to a calculated background of 30 electromagnetic tridents in the same region.

A. Rate Calculation

We have calculated the expected Υ rates using a photon-gluon-fusion (γGF) model[3] which accounts[4] for most of the published features[5] of ψ muoproduction. It uses a Bethe-Heitler diagram for heavy quark pair production with the nuclear photon replaced by a gluon. Additional soft gluon exchanges needed to conserve color are assumed not to affect the kinematics. The diagram is shown in figure V.1. Figures V.2 and V.3 illustrate the good agreement between the γGF model and ψ production by muons and photons. Using a distribution $G(x)=3(1-x)^5/x$ in gluon momentum fraction x, a bottom quark mass $m_b=4.7$ GeV/c^2, a bottom quark charge $|q_b|=1/3$, and a strong coupling constant $\alpha_s=-1.5/\ell n(4m^2_{b\bar{b}})$, where $m_{b\bar{b}}$ is the mass in GeV/c^2 of the produced quark pair, the model predicts Υ muoproduction cross sections of 0.13×10^{-39} cm^2 at 209 GeV and 0.28×10^{-39} cm^2 at 275 GeV. With $B(\Upsilon\rightarrow\mu^+\mu^-) = 3.1\pm0.9$ percent[5], the expected values of $B\sigma$ are $(4.0\pm1.2)\times10^{-39}$ and $(8.7\pm2.5)\times10^{-39}$ cm^2, respectively. The BCDMS upper limit is (70±40)% of the latter cross section.

B. Dimuon Mass Spectrum

Figure V.4 displays the spectrum in dimuon mass $M_{\mu^+\mu^-}$ from this experiment. Events below 5 GeV/c^2 in $M_{\mu^+\mu^-}$ are reconstructed and momentum fit as previously described. Above 5 GeV/c^2, the analysis of all events was checked by a hand reconstruction which was blind to the invariant mass. At all masses the assignment of beam-sign secondary muons either to the scattered muon or to the produced muon pair is the critical decision in the analysis. Incorrect pairing of muons from ψ or muon trident production can cause events which properly belong in the low-mass region to be misinterpreted as having a higher mass. Our muon pairing algorithm was selected primarily to minimize this problem. The scattered

muon is chosen to be the one with the smaller value of the square of its scattering angle divided by its scattered energy[6]. The algorithm is 89% efficient in reconstructing Υ's generated by the Monte Carlo simulation described below. The alternative choice for the scattered muon would produce more than a one-order-of-magnitude exaggeration of the high-mass continuum near the Υ, as shown by the "mispaired" histogram segment in figure V.4. We emphasize that the muon pairing algorithm can be optimized only if all three final-state muons are momentum-analyzed.

Despite the care exercised in muon pairing, Monte Carlo studies show that there remains a significant contribution in the region $4.7<M_{\mu^+\mu^-}<8.4$ GeV/c^2 from incorrectly analyzed lower-mass events. Allowance for these effects is most reliably made by use of an empirical fit to the mass continuum. This mass region, together with the range $1.5<M_{\mu^+\mu^-}<2.3$ GeV/c^2, was chosen for the fit in order to exclude regions complicated by charmonium production or rapid variations in low-mass acceptance. After subtraction of the fit continuum, the ψ peak in figure V.4 exhibits an 8.5% rms resolution, $\cong$1% larger than the Monte Carlo prediction[4]. The extrapolated continuum contains 1.8±1.0 background events in the Υ region $8.4<M_{\mu^+\mu^-}<11.1$ GeV/c^2, which in fact includes two observed events. The additional event at 11.5 GeV/c^2 is interpreted as continuum background with 65% probability, or as part of the peak corresponding to known Υ states with 1% probability. With 90% confidence, there are fewer than 3.8 events above the extrapolated background.

C. Acceptance Modeling

The Monte Carlo program used to simulate Υ muoproduction is based

on a routine which successfully parameterizes our ψ data[7]. There is a threshold in Q^2 for the virtual photoproduction of vector mesons by muons because the muon cannot lose energy and still stay on its mass shell without the photon acquiring some virtual mass. In order to reproduce the experimental ratio of coherent to incoherent ψ production from Fe nuclei, to parameterize threshold effects, and to describe the dependence on -t, the square of the four-momentum transferred to the target, the cross section is assumed to be

$$d\sigma/dt(\gamma Fe\rightarrow\psi X) = G(t)d\sigma/dt(\gamma N\rightarrow\psi N) \quad (t=0),$$

$$G(t) = A_e^2 \exp(\alpha t) + A_e\left[(1-\varepsilon\delta)\exp(\beta t) + \varepsilon\delta\exp(\delta t)\right].$$

The t resolution of the spectrometer is such that a δ-function at t=0 is smeared into $\sim\exp(5t)$. Therefore, data from other photon nucleus experiments[8] are averaged to set the coherent slope α to $150(\mathrm{GeV/c})^{-2}$. The shadowing factor A_e is taken to be $0.9x(A=55.85)$ based on electron-nucleus scattering data[9] at similar average Q^2. We have used $\beta=3$ $(\mathrm{GeV/c})^{-2}$, $\delta=1$ $(\mathrm{GeV/c})^{-2}$ and $\varepsilon=1/8$. These choices are consistent with high energy ψ photoproduction[10] and our experimental t distribution.

The ψ Monte Carlo is adapted to Υ simulation by appropriately scaling the vector-meson-mass-dependent parameters. Simulated Υ mass resolution and detection efficiency are 9% (rms) and 22%, respectively. The corresponding values for ψ production are 8.5%(rms) and 19%, showing the uniformity of the experiment over a wide range of dimuon invariant mass. The Υ cross section is normalized to the γGF value described above. Υ, Υ', and Υ'' states are generated in the ratio 1:0.39:0.32 in agreement with recent measurements of $\Gamma_{ee}(\Upsilon):\Gamma_{ee}(\Upsilon'):\Gamma_{ee}(\Upsilon'')$[11]. Υ' and Υ'' production suffer an additional $\sim$30% suppression relative to Υ production

because of threshold kinematics. The reconstructed peak corresponding to $10^4\times$ the expected signal is shown in figure V.4; 1.0 events from all Υ states are expected in the data.

D. Results

Our 3.8-event limit, integrated luminosity, and detection efficiency combine to set the 90%-confidence limit[12] $\sigma(\mu N\rightarrow\mu TX)B(\Upsilon\rightarrow\mu^+\mu^-)<22\times10^{-39}$ cm^2. With $B(\Upsilon\rightarrow\mu^+\mu^-)=(3.1\pm0.9)\%$[5], we obtain the 90%-confidence cross-section limit $\sigma(\mu N\rightarrow\mu TX)<0.79\times10^{-36}$ cm^2, including the error in the branching ratio. This limit lies above published predictions which use either the vector-meson dominance[13,14] or the γGF[15] models. Ignoring any γGF model uncertainty, this result rules out the choice $|q_b|=2/3$ with 85% confidence. With 67% confidence, the data disfavor the existence of similar bound states of a second charge 1/3 quark in the Υ mass region.

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Chapter V

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13. G. Aubrecht and W. Wada, Phys. Rev. Lett. 39, 978 (1977), make photoproduction predictions which, with VDF ν- and Q^2-dependence assumptions, correspond to 0.15×10^{-36} cm^2 of Υ muoproduction at 209 GeV.

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15. H. Fritzsch and K. Streng, Phys. Lett. 72B, 385 (1978), make photoproduction predictions which, with γGF ν- and Q^2-dependence assumptions, correspond to 0.36×10^{-36} cm^2 of τ muoproduction at 209 GeV.

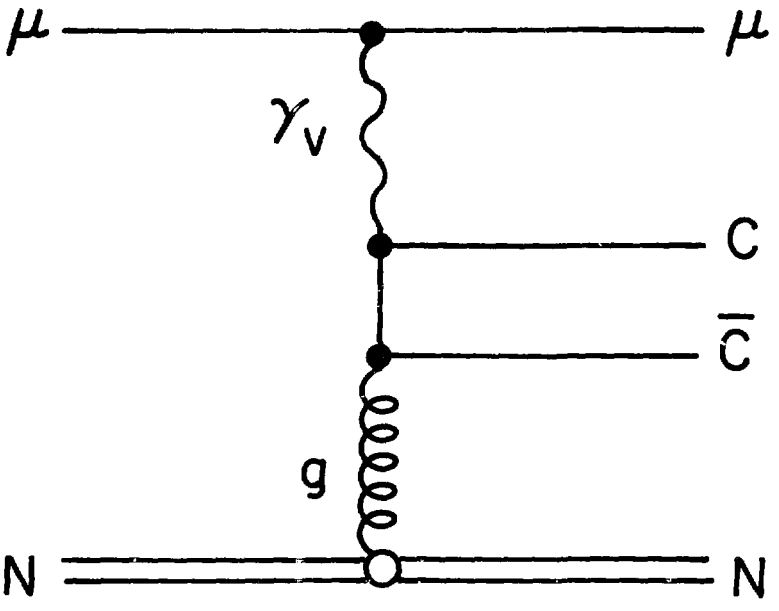
Figure Captions

Figure V.1. Feynman diagram for virtual photon-gluon-fusion production of charm states.

Figure V.2. Theoretical curve corresponding to the photon gluon cross section compared to $\gamma N \rightarrow \psi N$ data from this experiment (Muoproduction data) and from Ref. 10 (Photoproduction data). Figure from Ref. 3.

Figure V.3. Theoretical curve corresponding to the photon gluon cross section compared to ψ muoproduction data from this experiment. Figure from Ref. 3.

Figure V.4. Spectrum of 102 678 dimuon masses from 75% of the trimuon data. The background is fit by $\exp(a+bm+cm^2)$ in the regions of the solid curve with a χ^2 of 13.7 for 14 degrees of freedom, and is extrapolated along the dotted curve. The "mispaired" histogram segment illustrates the appearance of the mass spectrum if the alternative muon-pairing choice is made. The background-subtracted peak is shown in the lower corner; the expected peak from $10^4\times$ the Monte-Carlo simulated Υ, Υ', and Υ'' sample is shown in the upper corner, with the contribution from Υ' and Υ'' in black.



XBL 814-2258

Figure V.1

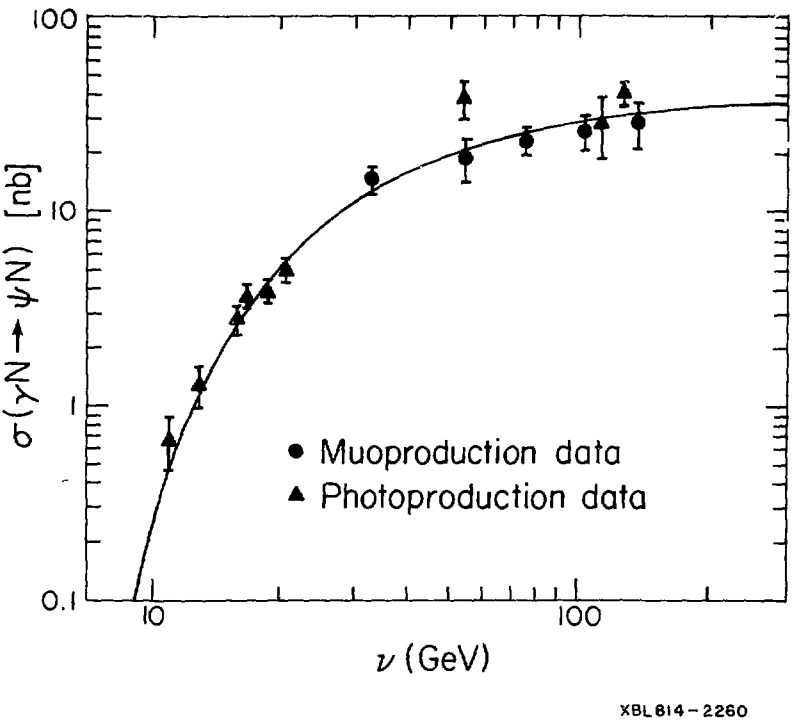

100
10
1
0.1
σ(γN → ψN) [nb]
10
100
ν (GeV)
● Muoproduction data
▲ Photoproduction data
XBL 814-2260

Figure V.2

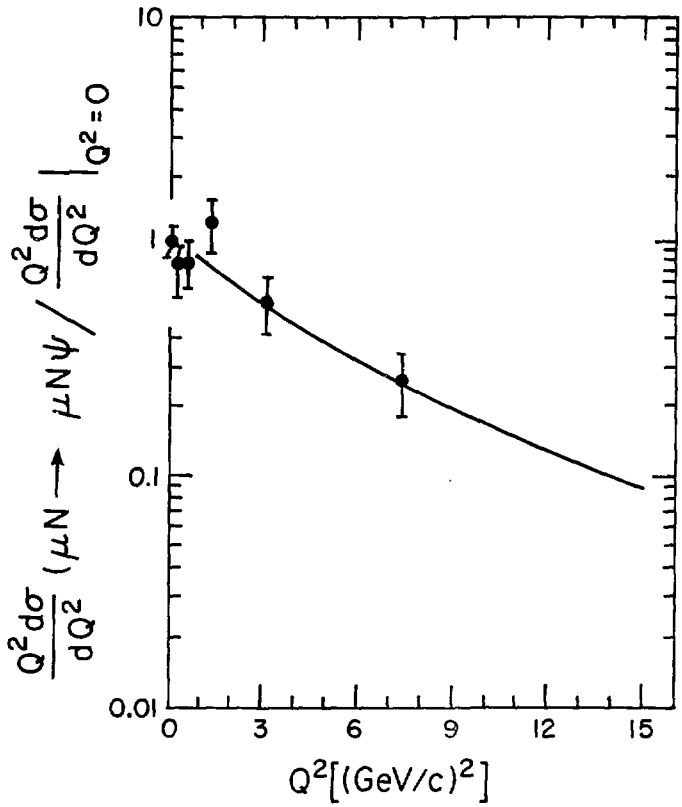
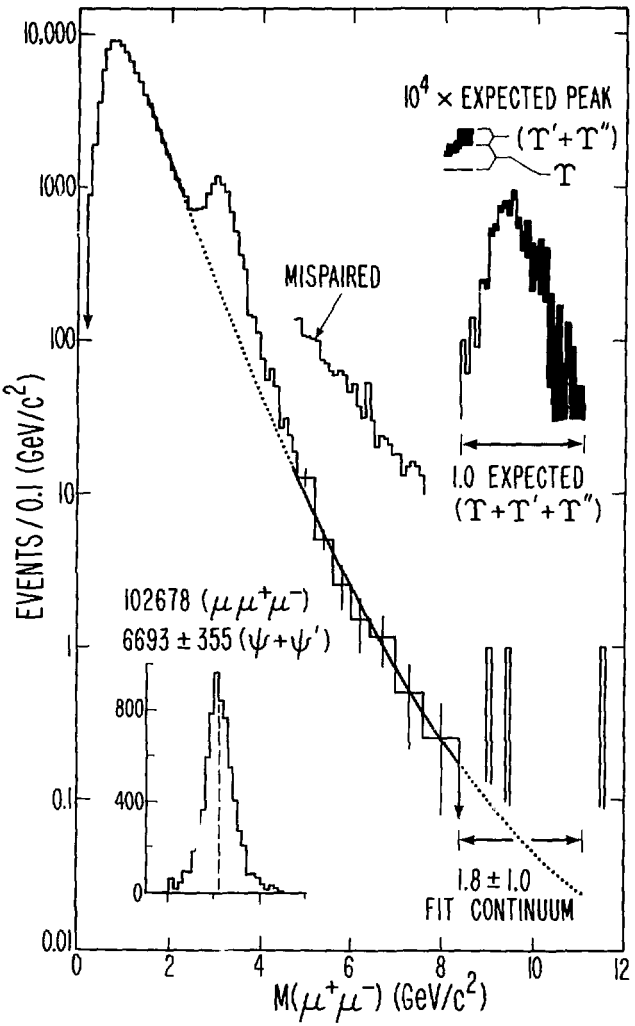

10
1
0.1
0.01
0
3
6
9
12
15
$Q^2[(GeV/c)^2]$
$\frac{Q^2 d\sigma}{dQ^2}(\mu N \rightarrow \mu N\psi) / \frac{Q^2 d\sigma}{dQ^2}\Big|_{Q^2=0}$
XBL814-2259

Figure V.3



XBL 805-966

Figure V.4

VI. LIMIT ON BOTTOM HADRON PRODUCTION

We have examined 36 952 dimuon final states produced by 1.4×10^{11} positive and 2.9×10^{10} negative 209-Gev muons. The majority of this data is due to the muoproduction of charmed hadrons, kaons, and pions, accompanied by their muonic decays. However, it is reasonable to enquire if there might be some contribution to this data from the muoproduction of hadrons containing bottom quarks with the subsequent muonic decay of these hadrons to charmed particles.

We have calculated the expected rate for bottom meson production using a photon-gluon-fusion (γGF) model, described previously, which accounts for most of the published features[1] of charmed meson production. Using, as before, a distribution $g(x)=3(1-x)^5/x$ in gluon momentum fraction x, a bottom quark mass $m_b=4.7$ GeV/c^2 and charge $|q_b|=1/3$, and a strong coupling constant $\alpha_s=1.5/\ell n(4m^2_{b\bar{b}})$, where $m_{b\bar{b}}$ is the mass of the produced quark pair, the model predicts a $b\bar{b}$ muoproduction cross section of 0.93×10^{-36} cm^2 at 209 GeV. If the $b\bar{b} \rightarrow \mu X$ branching ratio B is assumed to be 0.17 (essentially the same as that for $c\bar{c} \rightarrow \mu X$), the predicted σB is 0.16×10^{-36} cm^2.

A. Monte Carlo Calculations

Monte Carlo charm events were simulated by using the γGF model with a charmed quark mass of 1.5 GeV/c^2 and charge $|q_c|= 2/3$. For incoherent events, the same dependence on -t, not predicted by the model, was used as for the ψ analysis. Similarly, the same nuclear parameters were used for coherent events. Quark pairs carrying the full photon energy were transformed to D mesons using a fragmentation function[2] $D(z) = (1-z)^{0.4}$

where z is $2\bar{E}_D/m_{c\bar{c}}$ and $\bar{E}_D$ is the D energy in the $c\bar{c}$ rest frame. Charged and neutral D's were produced in a 1:2 ratio[2] and decayed to muons[3] with 20% and 4% branching ratios respectively[4]. Production and decay of other charmed states was not explicitly simulated. The $K\mu\nu$ ($K^*\mu\nu$) branching ratio was taken as 0.61 (0.39)[4]. The trigger efficiency for νGF charm events with decay muons is 16.7%.

Dimuon events from the decay in flight of muoproduced pions and kaons were simulated with a Monte Carlo using inelastic structure functions parameterized by the Chicago-Harvard-Illinois-Oxford collaboration[5]. The same experiment provided[6] the π and K production data used to determine final state particle multiplicities and momentum distributions. Bubble chamber data[7] was used to parameterize secondary meson-nucleon interactions. This use of experimental input made the Monte Carlo independent of models of hadron production. Hadron trajectories were simulated in the same detail as muon trajectories. The systematic normalization uncertainty in this Monte Carlo was determined to be ±50% by comparing the calculated π, K fraction with that obtained by representing the data as a combination of simulated π, K decay and charm events. The combined trigger and reconstruction efficiency for an event where a muon scatters and produces a muon from a π or K decay in the shower with an energy greater than 5 Gev is 4.6%.

Cuts are applied to reduce the contribution from π and K decay to (27±14)% of the dimuon sample. These cuts require a 9 GeV minimum daughter muon energy, a minimum ν of 75 GeV, a 0.2 GeV/c minimum daughter muon momentum, $p_\perp$, transverse to the virtual photon, and a range in inelasticity, $y=1-(\text{daughter muon energy})/\nu$, of $0.675<y<0.95$.

Histograms of simulated π- and K-decay events are subtracted bin by bin from the data histograms. Almost all of the remaining events are attributed to charmed meson decay. When these events are simulated with the γGF model, using the Monte Carlo program described above, background-subtracted data and charm Monte Carlo agree adequately in ν, Q^2, y, and daughter muon energy, while $p_\perp$ is higher in the data by 15%[1]. The measured cross section for diffractive charm production by 209 GeV muons is $6.9^{+1.9}_{-1.4}$nb.

Monte Carlo simulation of $b\bar{b}$ muoproduction is also based on the γGF model described above. As in the case of charm production, quark pairs carrying the full photon energy are transformed to B mesons using the fragmentation function $D(z) = (1-z)^{0.4}$, z is $2\bar{E}_B/m_{b\bar{b}}$, where $\bar{E}_B$ is the B energy in the $b\bar{b}$ reference frame. The B mesons decay to muons via $B \to D\mu\nu$. Further muon-producing cascade decays are ignored, because they tend to produce decay muons which are indistinguishable from charm background. The diffractive and shadowing parameters used are the same as those used in the ψ Monte Carlo. The simulated detection efficiency for $b\bar{b}$ states decaying directly to at least one muon is 19%.

B. Analysis Procedure

The ratio of simulated bottom quark events to simulated charm quark events is highest in the region ν>150 GeV and $p_\perp$>1.4 GeV/c. Hereafter we refer to this region as $R_{b\bar{b}}$. That $R_{b\bar{b}}$ should contain a higher ratio of $b\bar{b}$ to $c\bar{c}$ may be understood from a model independent viewpoint in that it takes a higher ν to create a heavier quark and a heavier quark produces more $p_\perp$ when it decays. The intent of the $b\bar{b}$ analysis reported here is to reshape slightly the $c\bar{c}$ Monte Carlo distributions in Q^2, y,

$p_\perp$, and ν in order to achieve full agreement with the data outside $R_{b\bar{b}}$. This procedure accounts for any inadequacies in modeling the data and reduces the dependence of this analysis on any particular model of heavy quark production. The reshaping is verified by requiring agreement between data and Monte Carlo in all kinematic spectra after all reshaping is completed. The empirically determined event-weighting functions which accomplish this reshaping are extrapolated into $R_{b\bar{b}}$, and are used to reshape the $c\bar{c}$ Monte Carlo distributions within that region. Since 58% of the events in $R_{b\bar{b}}$ have $\nu<170$ GeV and 50% have $p_\perp<1.6$ GeV/c the extrapolation is small for the majority of the events because the extrapolation covers a range which is only 27% of the kinematic range of the data on which it is based in ν and 17% of the range on which it is based in $p_\perp$. Furthermore, the extrapolation is done simultaneously in 2 dimensions in the $p_\perp$-ν plane, based on statistics 61 times those in $R_{b\bar{b}}$. The errors in the extrapolation are fully propagated and are included in all calculations. The spectra inside $R_{b\bar{b}}$ of the reshaped charm Monte Carlo and the background-subtracted data are compared to search for a possible $b\bar{b}$ signal.

The charm Monte Carlo spectra are reshaped by weighting each simulated $c\bar{c}$ event by a product of three functions, respectively of Q^2, y, and (ν and $p_\perp$). The weighting functions were $(1+Q^2/70(\mathrm{GeV/c})^2)^{-2}$, a polynomial[8] in y and the function of ν and $p_\perp$ listed in Table 1. The last function was determined by a two-dimensional fit in the ν-$p_\perp$ plane. Since Q^2 and y are only weakly correlated with $p_\perp$ and ν it was possible to determine the three weighting functions by iteration. After weighting by all three functions, each event was added to each histogram to produce the reshaped spectra. Before and after weighting, the charm

Monte Carlo sample was normalized to the background-subtracted data outside $R_{b\bar{b}}$.

C. Results

Figures VI.1 and VI.2 show background-subtracted data compared to the original and weighted $c\bar{c}$ Monte Carlo spectra in Q^2 and y. Also shown is 100× the $b\bar{b}$ signal (with $\sigma B = 0.16 \times 10^{-36}$ cm^2) expected from the γGF model. These spectra are populated only by events outside of $R_{b\bar{b}}$. Figures VI.3 and VI.4 make the same data-$c\bar{c}$-$b\bar{b}$ comparison. Figure VI.3 displays the ν spectra for $p_\perp > 1.4$ GeV/c and $p_\perp < 1.4$ GeV/c, and figure VI.4 shows the $p_\perp$ spectra for $\nu > 150$ GeV and $\nu < 150$ GeV. These figures emphasize the consistency between data and reshaped charm Monte Carlo outside $R_{b\bar{b}}$. Specifically, in the ν-$p_\perp$ plane outside $R_{b\bar{b}}$ the χ^2 for a unit ratio of data to $c\bar{c}$ Monte Carlo is 190 for 176 degrees of freedom.

The region $R_{b\bar{b}}$ contains 3.4 simulated $b\bar{b}$ events, or 29.5% of the Monte Carlo $b\bar{b}$ sample, and 455 $c\bar{c}$ events, or only 1.5% of the weighted Monte Carlo $c\bar{c}$ sample. After subtraction of the four simulated π- and K-decay background events, 456 data events remain in $R_{b\bar{b}}$. The error in the difference between data and Monte Carlo is $(\sigma_1^2+\sigma_2^2+\sigma_3^2)^{\frac{1}{2}}$, where $\sigma_1 = 22$ is the random error in the number of background-subtracted data events in $R_{b\bar{b}}$ and $\sigma_2 = 37$ is the error in the number of $c\bar{c}$ Monte Carlo events in $R_{b\bar{b}}$. Included in σ_2 are the random error in the ratio of Monte Carlo to data outside $R_{b\bar{b}}$, the error in weighting $c\bar{c}$ Monte Carlo events within $R_{b\bar{b}}$ based on the spectra outside $R_{b\bar{b}}$, and the random error in the generated number of these events. The error analyses which determine σ_1 and σ_2 take fully into account the statistical effects of variations in the amount of subtracted background and in the weights as-

signed to individual events. The systematic error induced by uncertainty in π- and K-decay background, σ_3=20, is determined by repeating the entire analysis with the background multiplied by 0.5 or 1.5. The resulting $b\bar{b}$ signal in the data is (1±48) events, corresponding to fewer than 62 candidates with 90% confidence. To ensure that any $b\bar{b}$ events outside $R_{b\bar{b}}$ do not affect the number of expected $c\bar{c}$ events in $R_{b\bar{b}}$, the analysis was repeated with 14× the simulated $b\bar{b}$ signal (corresponding to 48 events in $R_{b\bar{b}}$) added to the background-subtracted data. The simulated $c\bar{c}$ signal in $R_{b\bar{b}}$ changed by less than one event.

With our luminosity and calculated detection efficiency, these <62 candidates produce the 90%-confidence limit $\sigma(\mu N \to b\bar{b}X)B(b\bar{b}\to\mu X)<2.9\times10^{-36}$ cm^2. Using B=0.17, $\sigma(\mu N \to b\bar{b}X)<17\times10^{-36}$ cm^2. After factoring out the equivalent flux[9] of transversely polarized virtual photons, the muoproduction limit restricts $\sigma(\gamma N\to b\bar{b}X)<4.3$ nb at an average virtual photon energy of 160 GeV, when the same branching ratio assumption is made.

Our limits are greater than some published predictions using γGF calculations, but conflict with others and with several vector meson dominance (VMD) models. The γGF calculations in Refs. 10 and 11 predicted $\sigma(\mu N \to b\bar{b}X)=1\text{-}3\times10^{-36}$ cm^2 and 4×10^{-36} cm^2, respectively. Ref. 12 used a γGF model to derive $\sigma(\gamma N\to b\bar{b}X)=16$ nb at 160 GeV. The authors of Ref. 13 employed a γGF approach with a fixed strong coupling constant to get $\sigma(\gamma N\to b\bar{b}X)=0.2$ nb. They also obtained 0.02-0.05 nb with calculations using a running coupling constant with various gluon momentum distributions, but found 22 nb using VMD-based calculations. The VMD-model calculation of Ref. 14 yielded $\sigma(\gamma N\to b\bar{b}X)=25$ nb; Ref. 15 predicted (1-10 nb) on the basis of empirical formulae and a sum rule derived by Shifman

et al.[16]. The generalized VMD calculation in Ref. 17 found that the $b\bar{b}$ photoproduction cross section could be as high as 125 nb.

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Chapter VI

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TABLE 1. Weighting function $R(\nu,p_\perp)$ for daughter muon momentum, $p_\perp$, transverse to the virtual photon and beam muon energy loss ν.

$f = \log_{10}(p_\perp)$

$R(\nu,f) = P(\nu,f)\cdot F(f)$

$P(\nu,f) = 1.43+a_0\nu+b_0 f+c_0\nu\cdot f+d_0\nu^2+e_0 f^2$

$F(f) = (L_1(f)+L_2(f))/(L_3(f)+L_4(f))$

$L_i(f) = (a_i+b_i f)/(|c_i-f|^{d_i}+e_i)\ (1\le i\le 4)$

i	a_i	b_i	c_i	d_i	e_i
0	-.0022	-.086	-.0021	-9.3×10^{-6}	-.57
1	181	165	-.17	2.1	0.04
2	-.032	0.031	0.29	5.7	2.8×10^{-5}
3	44	3.9	-.20	2.6	0.010
4	-.0045	0.0074	0.30	6.4	9.8×10^{-6}

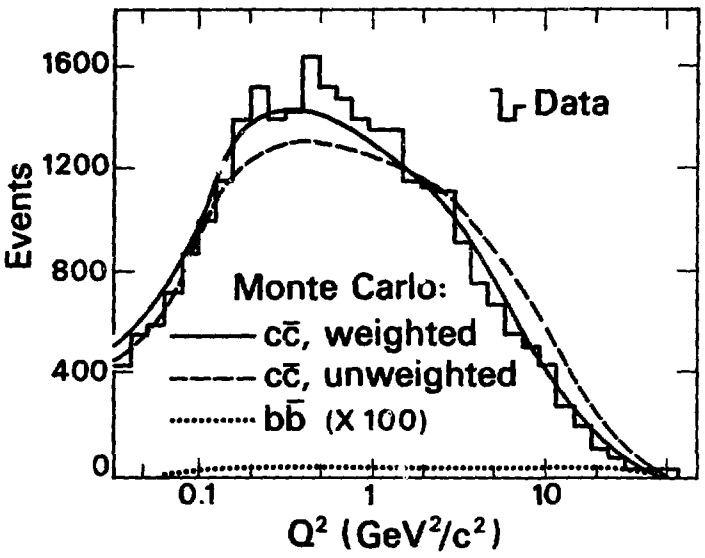
Figure Captions

Figure VI.1. Original and weighted $c\bar{c}$ Monte Carlo Q^2 spectra, compared with data after subtraction of the simulated π- and K-decay background. All events lie outside of $R_{b\bar{b}}$, the region where ν>150 GeV and the momentum, $p_\perp$, of the daughter muon transverse to the virtual photon exceeds 1.4 GeV/c. Also shown is the simulated Q^2 spectrum for 100× the $b\bar{b}$ signal expected from the γGF model.

Figure VI.2. Original and weighted $c\bar{c}$ Monte Carlo inelasticity y=1-(daughter muon energy)/ν, compared with background subtracted data, for events lying outside of $R_{b\bar{b}}$. Also shown is the simulated y spectrum for 100× the $b\bar{b}$ signal expected from the γGF model.

Figure VI.3. Original and weighted $c\bar{c}$ Monte Carlo ν spectra, compared with background subtracted data for (a) $p_\perp$>1.4 GeV/c and (b) $p_\perp$<1.4 GeV/c. Also shown are the simulated ν spectra for 100× the $b\bar{b}$ signal expected from the γGF model.

Figure VI.4. Original and weighted $c\bar{c}$ Monte Carlo $p_\perp$ spectra, compared with spectra of background subtracted data for (a)ν > 150 GeV and (b) ν<150 GeV. Also shown are the simulated $p_\perp$ spectra for 100× the $b\bar{b}$ signal expected from the γGF model.



XBL 8011-2411

Figure VI.1

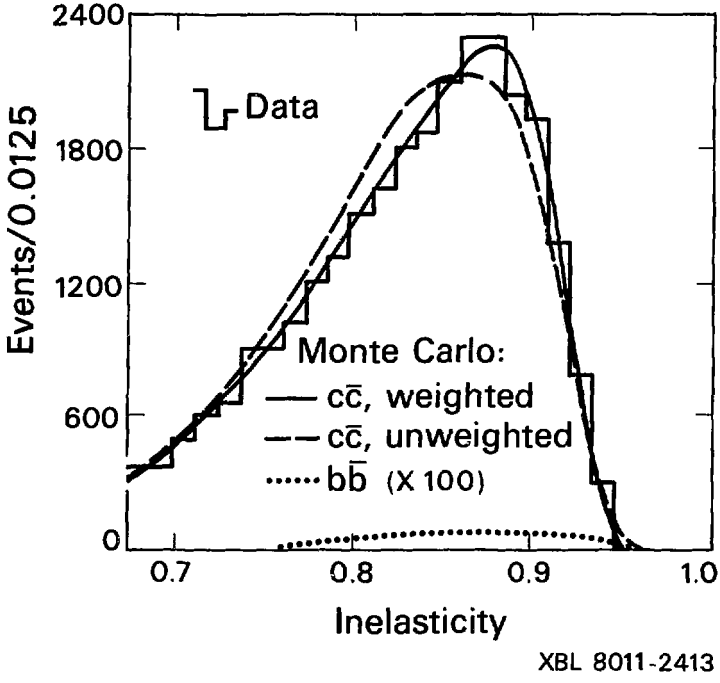

2400
1800
1200
600
0
Events/0.0125
Data
Monte Carlo:
— cc̄, weighted
– – cc̄, unweighted
····· bb̄ (X 100)
0.7
0.8
0.9
1.0
Inelasticity
XBL 8011-2413

Figure VI.2

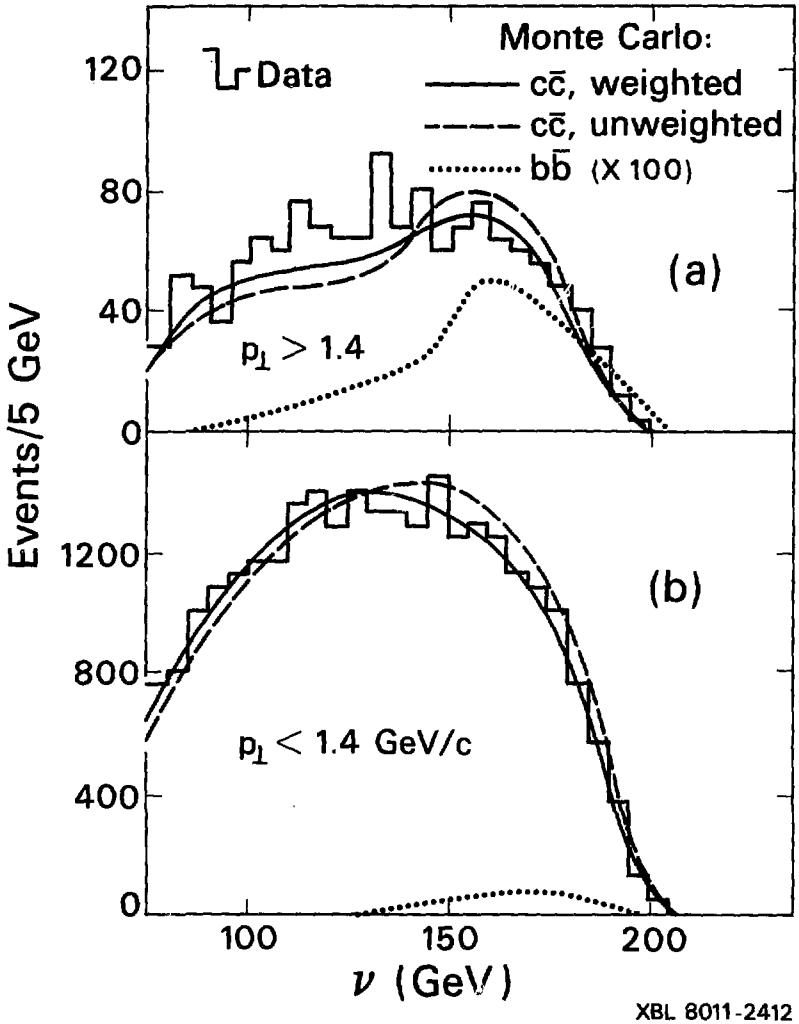

Monte Carlo:
Data
cc̄, weighted
cc̄, unweighted
bb̄ (X 100)
(a)
p⊥ > 1.4
(b)
p⊥ < 1.4 GeV/c
Events/5 GeV
ν (GeV)
XBL 8011-2412

Figure VI.3

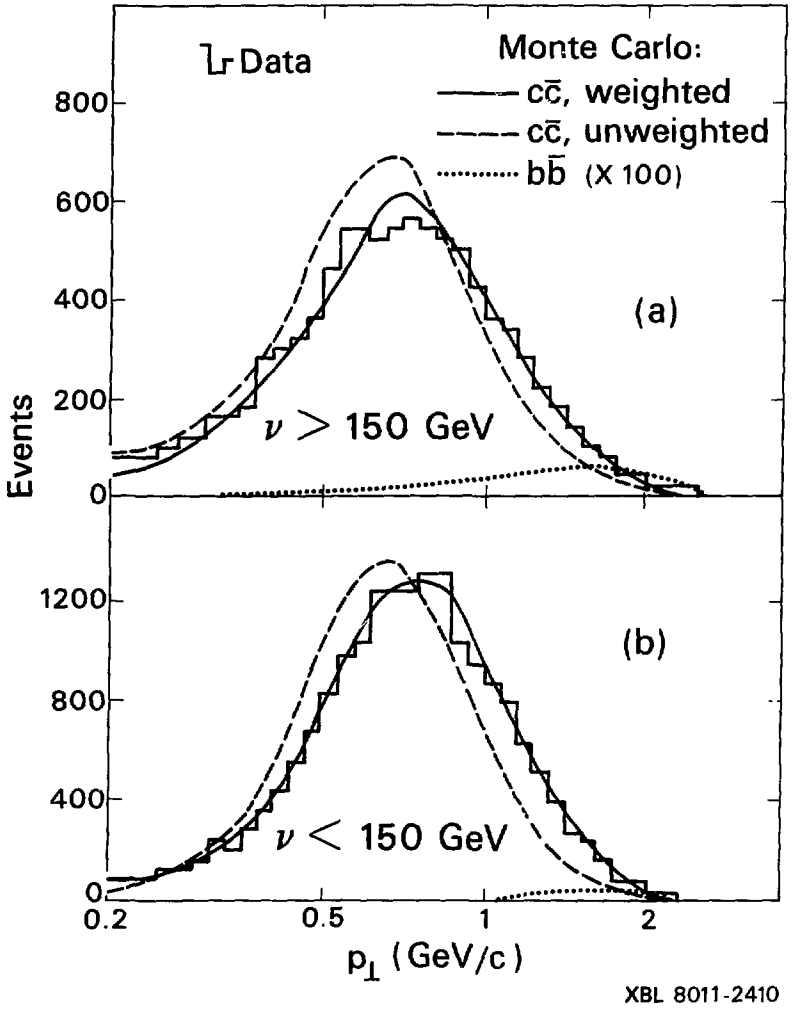

Data
Monte Carlo:
cc̄, weighted
cc̄, unweighted
bb̄ (X 100)
(a)
ν > 150 GeV
(b)
ν < 150 GeV
Events
800
600
400
200
0
1200
800
400
0
0.2
0.5
1
2
p⊥ (GeV/c)
XBL 8011-2410

Figure VI.4

VII. RARE MULTIMUON FINAL STATES

The large target and uniform acceptance of this experiment render it ideally suited for a search for small cross section processes that yield unusual numbers or topologies of muons in the final state. Two complete scans of events selected from the experimental sample produced by 1.4×10^{11} positive and 2.9×10^{10} negative 209-Gev muons have revealed sixteen 4-muon events and twelve 5-muon events. The integrated luminosity of 0.78×10^{39} cm^{-2} also produced 31 events of the type $\mu^{\pm}N \rightarrow \mu^{\pm}\mu^{\mp}\mu^{\mp}X$ and 13 events of the type $\mu^{\pm}N \rightarrow \mu^{\pm}\mu^{\pm}\mu^{\pm}X$. We refer to these two types as odd-signed trimuons to distinguish them from common trimuon production: $\mu^{\pm}N \rightarrow \mu^{\pm}\mu^{\pm}\mu^{\mp}X$. In every event all outgoing muons are fully momentum analyzed and their momenta are checked for energy conservation by including measurement of the incident muon momentum and calorimetric measurement of the associated shower energy. No reports of muon induced odd-signed trimuons or 4- or 5-muon final states have been published. Therefore we define these types of events as "rare" events.

A. Analysis

This sample of rare multimuon final states was culled from an initial sample of events in which the preliminary track reconstruction found sufficient candidate tracks which could be attached to the event vertex and provide the appropriate final state configuration of a rare event. Computer-drawn pictures of these events were scanned by physicists and the legitimate events were selected, for which $\sim$1 m^2 pictures were generated containing all raw wire chamber hits resolved to better than 1 mm in real transverse coordinates. With the high-resolution pictures, raw chamber hits are reconstructed by hand into tracks and the

vertex position is determined. The track reconstruction is then forced to fit the event using the hand-selected information. The information from the computer reconstruction as to the chi-square of each track and the probability that each chosen wire hit belongs on the track is examined and, if necessary, tracks are altered until the optimum event reconstruction is obtained.

To be accepted as a rare event, the result of the computer-assisted hand-forced fit is required to display the same topology as that of the original reconstruction. Close inspection of each high resolution picture insures that additional tracks crossing as few as 3 chambers have not been missed and that distinct tracks separated along their length by as little as 5 mm have not been combined. Figures VII.1, VII.2, and VII.3 show respresentative pictures of an odd-signed trimuon, a 4-muon event and a 5-muon event, respectively.

Several precautions assure that events are legitimate and ensure that two interactions are not mistakenly superimposed: The trigger demands only one beam track within a 57 nsec window centered on the event. All tracks are required to emanate from a tightly defined common vertex. All tracks are required to intersect the appropriate fine-grained hodoscope scintillators, sensitive within a ±10 nsec window. Adjacent drift and proportional chamber hits are required to register at a level rejecting tracks out of time by more than $\sim$50 nsec. The accepted tracks satisfy a tight x^2 cut separately in both orthogonal views. At least 3 hits in the third view link the two projections. Each accepted track, passing smoothly through >12 absorption lengths of steel can be interpreted only as a muon. The sign of each muon's charge is at

least 8 standard deviations from the reversed value.

Tables 1-6 present the properties of the rare multimuon events found by 2 complete scans of the data sample. These scans reveal scanning efficiencies of ∿90% for all three types of rare events. Of the original sample of events found by the scan and passing reconstruction, the following pass the tight cuts: 7 of the type $\mu^+N\to\mu^+\mu^+\mu^+X$, 22 of $\mu^+N\to\mu^+\mu^-\mu^-X$, 1 of $\mu^-N\to\mu^-\mu^-\mu^-X$, 6 of $\mu^-N\to\mu^-\mu^+\mu^+X$, 6 of $\mu^+N\to\mu^+\mu^+\mu^-\mu^-X$, 8 of $\mu^+N\to\mu^+\mu^+\mu^+\mu^-X$, 1 of $\mu^-N\to\mu^-\mu^-\mu^+\mu^+X$, 5 of $\mu^+N\to\mu^+\mu^+\mu^+\mu^-\mu^-X$, and 5 of $\mu^-N\to\mu^-\mu^-\mu^-\mu^+\mu^+X$. These events are produced in a data sample that contains 75,906 $\mu^\pm N\to\mu^\pm\mu^\pm X$, 112,369 $\mu^\pm N\to\mu^\pm\mu^\mp X$, and 110,626 $\mu^\pm N\to\mu^\pm\mu^\pm\mu^\mp X$. All of the events mentioned pass the same analysis cuts and all samples contain contributions from the "two muon" and "three muon" triggers.

B. Odd-Signed Trimuons

An intriguing possible cause for the odd-signed trimuons is a bottom hadron cascade, such as : $\mu^+N\to\mu^+b\bar{b}$; $\bar{b}\to\bar{c}\mu^+\bar{\nu}_\mu$, $\bar{c}\to$hadrons; and $b\to c+$hadrons, $c\to\mu^+\bar{\nu}_\mu+$hadrons. However, the limit on bottom hadron muoproduction set previously, when the muonic branching ratios and reconstruction efficiency are included, implies a maximum of 3 events from this source. The most probable cause of the odd-signed trimuon events is a dimuon produced by a charmed particle decay in which an extra muon from a π or K decay was produced in the hadronic shower. If the muon is of the correct sign, it will yield the final state muon charge configuration of an odd-sign trimuon. The charm dimuon signal is isolated from the data by subtracting off the absolutely normalized amount of π- and K-decay events from the entire dimuon sample[1]. The remaining 100,446

dimuons, which pass the same analysis cuts as 36 of the odd-signed trimuons, are ascribed to charmed particle decay. Besides the track χ^2 requirement these cuts also demand a shower energy greater than 12.5 GeV and an energy transfer ν greater than 30 GeV.

The expected number of odd-signed trimuons due to muoproduction in the shower of a charm dimuon may be estimated in two ways. Firstly, convolution of the shower energy spectrum of the charm dimuons with the Monte Carlo generated probability to obtain a muon from π or K decay versus shower energy[2] yields 70 events, of which 1/2, or 35 are expected to have the muon of the appropriate charge. Folding in the 50% uncertainty in the normalization of the π- and K-decay Monte Carlo produces the range 18-53 for this estimate. Secondly, one can observe directly the number of muons produced in showers of single muon inelastic scattering events. In a sample of 223,208 inelastic muon scattering events there are 146 events having a second muon with the opposite sign from that of the scattered muon, where this second muon can be attached to the event vertex and the event then passes analysis cuts. In all these events the second muon did not contribute to the event trigger. As an additional precaution against considering tracks that are not real, one can require events to have the total momentum of the outgoing track(s) not to exceed the incoming momentum by 52 GeV. This reduces the inelastic scatters to 222,158 and the oppositely-charged second-muon events to 132. This shows less than 9% of the 146 events, or 13 events are not real.

Of the 133 legitimate events, a certain number may be due to charmed particle decay. The measured charm muoproduction cross section

at 209 GeV times the branching ratio to muons is $1.29^{+.35}_{-.29}$ nb. Of the muons produced by charm, 64% exceed 5 GeV in energy. Therefore, the cross section to produ _ a muon with an energy greater than 5 GeV from a charm decay is .66-1.06 nb. The cross section to scatter and produce a muon from π or K decay with energy greater than 5 GeV is 2.28 nb. The muon from π or K decay has a 79% probability of being reconstructed, whereas the probability from a charm decay is 88%. This makes the ratio of the production of reconstructed muons with more than 5 GeV in energy from charm to that from π and K decay 0.32-0.52. Therefore, 65%-76% of the opposite sign second muon sample is due to muons produced by π or K decay in an hadronic shower. These 86-100 events yield the probability to produce a muon of a given charge in an hadronic shower of $(3.9\text{-}4.5)\times10^{-4}$. Therefore, we expect the charm dimuon sample to produce 39-45 odd-sign trimuons from hadronic shower muoproduction.

In order to further determine if the source of the odd-signed trimuons is hadronic shower muoproduction in the charm dimuon sample, figure VII.4 compares the kinematic spectra of the charm dimuon sample with those of the odd-signed trimuons. We apply a statistical test to these distributions to determine their mutual consistency. The Kolmogorov-Smirnov test is superior to the usual Pearson's χ^2 test for small samples and does not involve the binning of individual observations[3]. Given n independent observations of a variable X denoted X_i, numbered in order of increasing magnitude, define

$$S_n(X) = \begin{cases} 0; & X < X_1 \\ i/n; & X_i \leq X \leq X_{i+1} \\ 1; & X \geq X_n \end{cases}$$

then the Kolmogorov-Smirnov test consists of finding the maximum of the absolute value of the difference between the $S_n(X)$ for the two distributions. This maximum is then converted into a confidence level through use of calculated tables[4].

The Kolmogorov-Smirnov test shows some deviations between the dimuon and odd-sign trimuon samples because the generation of an additional muon in a charm dimuon affects the event topology so that the event will appear slightly altered from a typical charm dimuon event even when reconstructed by an analysis blind to the third muon. We believe this effect is probably most pronounced in assessing the inelasticity and shower energy of events. Table 7 presents the results of the Kolmogorov-Smirnov test on the two samples. They are adequately consistent. The comparison of the $p_\perp$ spectra is particularly important in that heavy quark production would produce a large inconsistency since the average bottom decay produces a $p_\perp$ of 1 GeV/c and the average charm decay 0.4 GeV/c. It also should be noted that the six spectra presented in figure VII.4 do not display independent variables if one assumes the parent process involves virtual photoproduction. However, the six could be less correlated were some other "new physics" involved in their creation.

C. Elastic 4- and 5-Muon Events

We observe three 4-muon events and five 5-muon events with a shower energy less than 6 GeV that pass our analysis quality cuts. We define these as elastic events. There are two 5-muon events not included in the elastic sample where the fifth muon track has a poor χ^2 and the remaining four tracks pass the χ^2 cut. The elastic 5-muon events are

probably due to electromagnetic tridents where an extra electromagnetic pair is produced and the 4-muon events are 5-muon events where the fifth, presumably low energy, muon was not seen.

The sources of electromagnetic pairs are shown in figure VII.5 for the case of electromagnetic trident production. We have done Monte Carlo studies of these processes and conclude that Bethe-Heitler dominates over bremstrahlung by a factor of 100. Since our experiment does not impose an opening angle cut on the outgoing muons, this ratio agrees with that found by Ref. 5 for a coherent iron target without cuts. We therefore believe the dominant contributions to the elastic 4- and 5-muon events to be the double Bethe-Heitler diagram shown in figure VII.6a.

In order to study the double Bethe-Heitler process we first consider single Bethe-Heitler events which constitute 99% of our electromagnetic trident sample. Examination of the elastic (shower energy greater than 6 GeV) trimuon sample reveals a large contribution from elastic psi production. The number of elastic psis is determined by fitting the dimuon invariant mass continuum above and below the region of charmonium production, extrapolating this fit into the region of charmonium production and subtracting the fit number of continuum events from the total in this region. The remaining events are ascribed to ψ and ψ' production. This number of elastic ψ and ψ' Monte Carlo events[6] is then subtracted from the entire elastic trimuon sample, leaving 87,650 events attributed to electromagnetic trident production. All of these events pass the same analysis quality cuts as the 4- and 5-muon events.

The expected number of elastic 4- and 5-muon events due to elec-

tromagnetic tridents generating an additional pair via a double Bethe-Heitler process is estimated two ways. Firstly, and most simply, these events are expected to appear with a frequency of $O(\alpha^2)$ less than electromagnetic tridents. This predicts 6 events. Secondly, the probability for a photon with sufficient energy to produce a muon pair, where each muon exceeds the detection threshold energy of 5 GeV, may be determined by comparing the total electromagnetic trident sample with the virtual photon flux that produced it. Inelastic ψ and ψ' events are subtracted off the inelastic trimuon sample as in the elastic case to determine the inelastic portion of the electromagnetic trident sample. When added to the elastic tridents, they comprise the total 104,496 events in the electromagnetic trident sample.

The equivalent flux[7] of transversely polarized virtual photons per muon is multiplied by the incoming flux of 1.7×10^{11} muons. The data corresponds to 2.04×10^{9} virtual photons with $\nu > 10$ GeV. This yields a probability of 5.1×10^{-5} to produce an extra pair, and have it trigger and be reconstructed. In the entire sample of 4- and 5-muon events 52% ± 19% would not have triggered without the presence of the additional muons beyond the spectator and the most energetic daughter muon of each sign. Therefore, folding in its additional probability for triggering, the expected rate for a virtual photon to produce an additional electromagnetic pair is $(1.1 \pm 0.37) \times 10^{-4}$. This then predicts 9.6 ± 3.2 elastic electromagnetic 4- and 5-muon events.

To test the hypothesis that the elastic 4- and 5-muon events are due to double Bethe-Heitler production, they may be compared with the events principally due to single Bethe-Heitler production, the elastic

electromagnetic tridents. Figure VII.7 compares the spectra of various kinematic quantities for the elastic 4- and 5-muon events with the elastic tridents. Table 8 presents the probability that these various kinematic spectra are consistent based on the application of the Kolmogorov-Smirnov test. The conclusion is that the elastic electromagnetic tridents form the parent sample of the elastic 4- and 5-muon events.

D. Inelastic 4-Muon Events

There are thirteen 4-muon events which have a shower energy greater than 6 GeV. Of these inelastic events there are 11 which have a shower energy greater than 12.5 GeV and a ν greater than 30 GeV. We believe these events are inelastic dimuons, primarily due to charm particle production with muonic decay, accompanied by the electromagnetic production of a muon pair. The diagram for this reaction is shown in figure VII.6b. After subtraction of the π- and K-decay background there are 100,446 dimuon events passing analysis cuts with a shower energy greater than 12 GeV, ν greater than 30 GeV. These are ascribed principally to charm meson production with a muonic decay. The previously determined probability to electromagnetically produce a muon pair of $(1.1\pm0.37)\times10^{-4}$ yields 11 ± 3.7 4-muon events expected from charm events with an additional electromagnetic pair.

Figure VII.8 compares the spectra of various kinematic quantities for the 4-muon events and the background subtracted dimuon events, where all events have a shower energy exceeding 12.5 GeV and a ν exceeding 30 GeV. Table 9 presents the probability that these spectra are consistent, based on the application of the Kolmogorov-Smirnov test. The

conclusion is that charm dimuons electromagnetically producing a muon pair are the most likely source for these inelastic 4-muon events.

Another possible source of the thirteen inelastic 4-muon events is that of an inelastic trimuon with an additional muon from a π or K decay in the hadronic shower. The inelastic (shower energy greater than 6 GeV) portion of the electromagnetic trident sample includes 16,845 events. The previously determined probability to produce a muon of a given charge in an hadronic shower exceeding 6 GeV of $(3.9\text{-}4.5)\times10^{-4}$ predicts 6-8 muons of each sign produced in the hadronic showers of the inelastic tridents. Thus as many as 12-16 of the 4-muon events may be expected from muoproductio n the hadronic showers of the inelastic tridents. The spectra of various kinematic quantities of the inelastic 4-muon events are ompared with the spectra for the inelastic tridents in figure VII.9.

Table 10a presents the probability that the spectra of the combined inelastic 4- and 5-muon sample are consistent with those of the inelastic tridents, and table 10b presents the probability that the spectra of the inelastic 5-muon events are consistent with those of the inelastic tridents. These probabilities, based on the Kolmogorov-Smirnov test, show that while the spectra of the combined sample are not consistent with the inelastic tridents, the inelastic 5-muon events by themselves are consistent. Therefore the source of the inconsistency between the combined sample and the inelastic tridents is due to the inelastic 4-muon events. It is evident that the contribution of inelastic tridents with hadronic shower muoproduction to the inelastic 4-muon sample must be small. The primary source of the inelastic 4-muon events is charm

production with electromagnetic pairs.

One inelastic 4-muon event bears further examination. Event 1191-5809 has an unusually high transverse momentum with respect to its virtual photon of 2.3 GeV. The probability that the two conventional processes here considered to be the source of the 4-muon events would produce one or more 4-muon events with a p_t greater than or equal to that of event 1191-5809 is 11%. The invariant masses of the two possible muon pair combinations are 3.5 and 3.0 GeV. The probability of producing an inelastic 4-muon event with a reconstructed invariant mass within one standard deviation (9%) of the ψ mass is also 11%. These and other considerations have prompted the interpretation of this event as diffractive $b\bar{b}$ production with $\bar{b}\to\psi X, \psi\to\mu^+\mu^- X$, and $b\to\mu^- \bar{\nu}_\mu X$[8].

E. Inelastic 5-Muon Events

There are five 5-muon events with a shower energy greater than 6 GeV. The most probable source for these events is that of an inelastic trimuon with an additional electromagnetically produced muon pair. The number of events due to such an inelastic double Bethe-Heitler process may be estimated by using the previously determined probability to electromagnetically produce a muon pair of $(1.1\pm0.37)\times10^{-4}$. This probability, when multiplied by the inelastic trident sample of 16,845 events yields 2 expected inelastic 5-muon events.

Another possible source of muon pairs would be the production in the hadronic shower of the inelastic tridents. However, the cross section for muon induced hadronic pair production in Ref. 9 is less by a factor 23 than the cross section for the muon induced Bethe-Heitler pro-

cess in Ref. 10. As mentioned earlier, other radiated sources of pairs are suppressed by a factor of 100 with respect to Bethe-Heitler. Figure VII.9 displays the spectra of various kinematic quantities of the inelastic 5-muon events with the spectra of the inelastic tridents. Table 10b presents the probability that the spectra of the inelastic 5-muon events are consistent with those of the inelastic tridents. The conclusion is that the inelastic 5-muon events appear due to the inelastic double Bethe-Heitler process, although their rate is unexpectedly high.

It is interesting to observe the sign of the beam muon producing the 5-muon events. The data sample which contains these events was induced by 1.4×10^{11} μ^+ and 2.9×10^{10} μ^-, a ratio of μ^+/μ^- of 5. However, of the five inelastic 5-muon events, three were produced by the μ^- beam. Overall, for the entire 5-muon sample, five are μ^+ induced and five are μ^- induced. One of the μ^- induced 5-muon events, 851-11418, has particularly remarkable characteristics in that it has a Q^2 of 3 GeV and a total transverse momentum with respect to the virtual photon of 2 GeV. The probability that the double Bethe-Heitler process would produce one or more events with a $p_\perp$ and Q^2 greater than or equal to the values of event 851-11418 is 3%.

F. Other Observations

Although there have been no other observation of muon induced rare multimuon events, there have been observations of neutrino induced odd-sign trimuons and 4-muon events. The CERN-Dortmund-Heidelberg-Saclay (CDHS) group reported[11] observing four $\nu \to \mu^- \mu^+ \mu^+$ with a calculated background of 6 events from π and K decays. They also observe[12] one event of the type $\bar{\nu} \to \mu^+ \mu^- \mu^-$. These events occur at a rate of 1×10^{-6} relative

to charged current neutrino scattering. The CDHS group has also observed[13] one event of the type $\nu \to \mu^+\mu^-\mu^+\mu^-$. The rate corresponding to the 4-muon event relative to the opposite sign neutrino induced dimuon events is $\sim 1.4 \times 10^{-4}$.

The Berkeley-Fermilab-Hawaii-Seattle-Wisconsin group has observed one event of the type $\bar{\nu} \to \mu^+e^-e^+e^-$ in the 15 foot bubble chamber at Fermilab[14]. The rate relative to single muon production for this event is of order 10^{-7}, the same as that corresponding to the CDHS 4-muon event. It is important to remember when comparing the muon and neutrino induced rare events that in the former case the model involves the interaction of a virtual photon with a sea charm quark and in the latter the interaction of a virtual W with a valence d or s quark.

The rare multimuon events reported here appear to be produced by conventional physics with the possible exception of one elastic 5-muon event and one inelastic 4-muon event. Nevertheless, diagrams such as those in figure VII.6 have not been observed before. The actual and expected numbers of events of all types are shown in table 11. To summarize, the odd sign trimuons have a rate relative to the dimuons of 3.6×10^{-4} and are due to charm dimuon events accompanied by an additional π or K decay. The elastic 4- and 5-muon events are electromagnetic in origin, specifically due to the double Bethe-Heitler process and have a rate relative to the elastic tridents of 9×10^{-5}.

The inelastic 4-muon events appear to be charm dimuons with an electromagnetically produced muon pair. There could also be a small contribution from inelastic tridents where a muonic π or K decay ocurred in the hadronic shower. The inelastic 4-muon events occur at a rate of

1.1×10^{-4} relative to the charm dimuons. The kinematics of the inelastic 5-muon events are consistent with those of inelastic tridents that electromagnetically produced an additional muon pair. However, their rate relative to the inelastic tridents is 3×10^{-4}, a rate that is higher than the 4-muon rate relative to the dimuons. This is anomalous because both types of event should display the same rate with respect to their parent process if both are due to electromagnetic pair production in their parent process. The observed rate of the 4-muon events with respect to the dimuons is consistent with the calcluated one, whereas the rate for the 5-muon events with respect to the inelastic tridents is not. This anomaly may suggest new physics when considered with the fact that although the ratio of incident positive to negative muon beam fluxes is 5:1, there is an equal number of 5-muon events induced by beam muons of each sign. However, the statistics are far from conclusive.

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Chapter VII

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Table 1. Odd sign trimuons listed by event number followed by the charge of the first through third outgoing muon and their momenta. Events are produced by an incident μ^+ beam except where noted.

EVENT	SIGNS	p_{1X}	p_{1Y}	p_{1Z}	p_{2X}	p_{2Y}	p_{2Z}	p_{3X}	p_{3Y}	p_{3Z}
533- 4135	+++	-1.8	0.6	41.1	-0.2	-0.1	29.5	0.8	-0.5	27.8
544- 284	+++	1.3	0.3	111.0	-0.5	-0.7	7.5	-0.4	0.3	6.7
555-11180	+--	1.6	0.2	143.6	-1.0	-0.3	12.5	1.3	0.2	8.9
588- 959	+--	-0.4	-0.3	102.4	-1.2	-0.8	20.7	0.2	-0.1	9.7
588- 1916	+--	-0.3	-0.2	37.3	-1.2	-0.4	10.1	-0.8	0.4	8.0
611- 3961	+--	-0.0	-0.0	23.5	0.9	-0.1	17.5	-0.4	0.2	14.3
643- 2708	+--	-1.3	0.1	79.9	-0.3	-0.7	19.7	0.3	0.3	8.7
644- 8059	+--	-0.4	-0.1	41.8	-0.4	0.7	23.0	-0.6	0.3	9.8
652- 6550	+--	-0.3	0.5	28.8	-1.0	-0.1	20.8	-0.2	0.2	9.3
666- 8769	+--	0.9	1.9	63.9	0.4	0	18.3	0.1	0.0	8.2
740- 2613	+--	0.9	0.2	86.5	0.0	0.1	20.0	-0.1	-0.2	11.5
770-10018	+--	1.3	1.5	78.2	1.0	-0.1	45.2	0.9	-1.2	11.0
773- 7250	+--	-1.5	1.4	53.4	0.4	-0.4	-24.8	-0.3	-0.3	9.1
808- 5590	+++	-0.3	-0.2	41.2	0.1	-0.8	25.4	-0.6	-0.4	8.8
830- 657	+--	-0.4	0.0	45.2	-0.4	-0.3	16.6	-0.1	0.3	9.4
847- 2596	-++*	0.3	-1.3	44.9	-1.0	0.3	30.9	0.0	0.4	23.2
847- 6635	-++*	-0.3	-0.5	86.4	-0.3	0.7	18.9	0.5	-0.1	10.4
851- 5726	-++*	-0.6	-0.1	48.9	-0.3	-0.4	18.6	-0.0	0.1	10.7
852- 9466	---*	-0.0	0.4	32.2	-0.8	-0.6	23.6	-0.0	-0.4	10.0
864- 3605	-++*	0.7	-0.0	98.5	-0.3	-0.2	12.0	0.1	-0.2	10.8
873- 7911	-++*	0.4	-0.2	34.5	0.9	-0.9	8.7	1.3	-0.0	8.1
885- 3661	-++*	-0.0	-0.4	45.8	0.6	-0.0	18.7	-0.2	0.6	12.4
928- 5026	+--	0.3	-0.1	101.2	0.0	-0.2	20.3	-0.6	-1.3	16.7
932-10333	+++	-0.3	-0.0	59.8	-0.2	0.0	20.9	-0.1	0.1	13.5
975- 7110	+--	-1.8	0.3	49.3	0.0	1.2	63.6	-0.0	-0.4	12.7
981- 1241	+++	-1.4	0.6	132.4	0.5	-0.4	15.5	0.3	-0.1	7.0
1001- 4560	+--	-2.1	-0.8	99.7	-0.2	0.8	11.2	-0.0	0.1	9.5
1010- 530	+++	0.6	-0.1	39.0	-0.1	-0.1	7.8	-0.3	-1.0	7.7
1013- 7037	+--	0.5	0.4	27.7	1.4	-0.3	45.3	0.4	-0.1	14.0
1028- 8809	+--	-1.0	0.0	85.6	0.3	-0.4	17.1	-0.8	-0.4	9.8
1035- 8075	+++	-0.6	0.1	102.8	0.0	-0.4	11.8	-0.0	-0.4	11.5
1057- 7403	+--	1.3	0.3	175.7	0.4	-0.7	9.3	-0.3	0.2	7.2
1118- 9435	+--	-0.8	0.8	98.6	0.0	-0.1	17.0	-0.5	-0.4	16.0
1132- 4519	+--	-0.3	0.3	67.7	-0.5	-0.2	8.6	-0.3	-0.2	7.3
1202- 9314	+--	-0.0	0.5	77.5	0.9	-0.3	14.0	0.3	0.1	10.5
1213- 940	+--	0.9	0.4	145.0	0.2	0.0	12.3	0.5	-0.3	8.0

*= μ^- beam

Table 2. Odd sign trimuons listed by event number followed by the shower energy deposited in the calorimeter, E_{shwr}, energy transfer ν, momentum transfer squared, Q^2, inelasticity Y, the momentum of the daughter muons, $p_\perp$, perpendicular to the virtual photon direction, and the missing energy, E_{miss}.

EVENT	E_{shwr}	ν	Q^2	Y	$p_\perp$	E_{miss}
533- 4135	112.0	174.0	18.49	.67	0.16	-3.2
544- 284	33.7	79.3	0.11	.82	1.06	31.2
555-11180	35.6	58.7	0.47	.63	0.19	1.5
588- 959	81.9	113.0	0.23	.73	1.37	0.6
588- 1916	154.2	176.6	0.19	.90	1.80	4.1
611- 3916	123.8	180.0	1.41	.82	0.11	24.4
643- 2708	93.0	132.4	0.74	.79	0.42	10.9
644- 8059	90.3	170.7	0.07	.81	1.25	47.6
652- 6550	138.1	183.6	1.55	.84	1.04	15.4
666- 8769	75.6	138.8	11.69	.81	0.51	36.6
740- 2613	58.8	122.4	0.01	.74	0.47	32.0
770-10018	69.4	128.7	5.81	.56	1.60	3.1
773- 7250	123.5	151.7	8.91	.78	0.71	-5.6
808- 5590	72.7	208.1	0.33	.84	1.30	101.2
830- 657	146.0	107.0	0.09	.84	0.20	-11.1
847- 2956	74.3	166.4	9.44	.68	0.63	38.1
847- 6635	73.9	124.2	0.73	.76	0.72	21.0
851- 5726	109.7	156.1	0.16	.81	0.33	17.1
852- 9466	119.5	183.6	1.18	.82	1.13	30.5
864- 3605	73.6	114.2	0.04	.80	0.50	17.7
873- 7911	123.4	173.7	0.20	.90	2.23	33.3
885- 3661	88.5	165.5	0.78	.81	0.78	46.0
928- 5026	51.8	138.9	0.86	.73	1.60	50.1
932-10333	136.0	151.5	0.04	.77	0.15	-18.8
975- 7110	37.9	130.6	2.61	.42	1.43	16.4
981- 1241	37.6	72.3	0.34	.69	1.19	12.3
1001- 4560	48.1	105.8	3.30	.80	0.65	37.0
1010- 530	169.3	163.6	0.22	.91	1.28	-21.3
1013- 7037	97.1	178.2	1.20	.67	1.03	21.7
1028- 8809	75.3	116.0	0.53	.78	0.98	14.3
1035- 8075	100.8	112.6	0.10	.79	0.78	-11.5
1037- 7403	16.4	36.4	0.26	.55	0.29	3.5
1118- 9435	53.4	109.5	1.79	.70	0.27	23.0
1132- 4519	168.5	140.5	0.50	.89	0.65	-43.9
1202- 9314	81.0	134.3	1.08	.82	0.95	28.8
1213- 940	22.3	50.1	0.27	.60	0.77	-2.4

Table 3. Four-muon events listed by event number followed by the charge of the first through fourth outgoing muon and their momenta. Events are produced by an incident μ^+ beam except where noted.

EVENT	SIGNS	p_{1X}	p_{1Y}	p_{1Z}	p_{2X}	p_{2Y}	p_{2Z}	p_{3X}	p_{3Y}	p_{3Z}	p_{4X}	p_{4Y}	p_{4Z}
538- 1662	++--	-0.5	0.1	20.7	-0.4	-0.0	13.5	-0.5	-0.5	20.4	0.1	-0.0	15.6
547- 7704	+-++	0.3	0.0	20.9	-0.0	0.2	19.2	-0.7	0.6	27.4	-0.5	0.1	13.5
550- 9806	++--	0.4	0.3	15.6	0.2	-0.3	6.6	0.4	-0.0	30.3	-0.2	-0.3	16.2
613- 3277	+-++	1.2	0.2	76.8	0.4	-0.1	16.5	1.7	0.7	17.0	0.0	0.0	8.9
672- 445	+-++	1.2	-0.0	96.3	0.1	-0.0	9.0	0.3	0.6	16.3	0.0	0.4	13.1
738- 4419	+-++	1.3	-0.2	100.5	1.3	-0.1	39.6	-0.1	-0.0	30.2	-0.5	0.1	8.9
777- 7592	++--	-1.9	-0.3	142.4	0.1	-0.3	19.0	-0.6	0.3	31.2	-0.5	-0.1	23.1
898- 1342	--++*	-1.1	0.5	86.9	-0.0	0.2	11.4	0.6	0.4	24.6	0.1	-0.2	9.3
1005- 3384	+-++	-2.7	0.3	175.0	-0.0	-0.3	28.7	-0.8	0.2	12.6	0.4	-0.6	9.4
1025- 6845	+-++	-2.1	-1.1	141.3	-0.7	0.8	57.5	0.3	0.2	16.1	-0.3	-1.0	11.6
1034- 3903	++--	-1.0	0.1	52.0	-1.1	-0.8	31.9	1.4	0.7	29.7	-0.1	-0.0	9.9
1079- 1845	+-++	-0.8	-1.0	58.2	-2.1	0.8	64.5	-0.6	-0.2	43.2	-1.0	0.0	11.7
1138-10327	+-++	-0.7	-0.1	176.9	0.0	-0.5	12.0	-0.2	0.5	13.4	0.2	-0.0	13.2
1141- 4818	++--	-0.4	-0.2	71.8	0.4	-0.0	27.9	-0.5	0.9	20.0	-0.4	0.1	10.5
1191- 5809	++--	-0.3	0.6	65.5	0.5	-2.2	25.5	1.9	1.0	28.8	0.0	1.0	17.8

*= μ^-beam

Table 4. Four-muon events listed by event number followed by the shower energy deposited in the calorimeter, E_{shwr}, the energy transfer ν, momentum transfer squared, Q^2, the momentum of the daughter muons, $p_\perp$, perpendicular to the virtual photon direction, the missing energy, E_{miss}, and the invariant masses formed by muons 2 and 3 and muons 2 and 4.

EVENT	E_{shwr}	ν	Q^2	$p_\perp$	E_{miss}	M_{23}	M_{24}
538- 1662	76.9	180.9	0.80	0.52	54.5	0.45	0.56
547- 7704	105.4	198.3	0.67	0.28	32.8	0.37	0.37
550- 9806	59.2	186.6	1.69	0.63	74.2	0.64	0.52
613- 3277	24.7	126.2	0.27	1.42	59.2	1.51	0.34
672- 445	39.1	99.7	0.03	0.95	22.2	0.52	0.48
738- 4419	48.1	110.1	0.96	0.74	-16.6	1.24	1.63
777- 7592	-4.0	62.4	0.26	0.62	-7.0	0.81	0.60
898- 1342	72.0	119.7	0.97	0.66	2.5	0.42	0.48
1005- 3384	3.9	24.9	0.46	0.92	-29.7	1.24	1.22
1025- 6845	6.3	76.3	2.17	1.15	-15.2	0.89	2.64
1034- 3903	77.8	154.9	0.37	1.13	5.5	2.92	0.69
1079- 1845	29.5	162.2	0.35	2.12	-1.9	1.43	1.42
1138-10327	0.3	34.8	0.01	0.13	-4.0	1.03	0.57
1141- 4818	48.2	146.2	0.31	0.84	39.6	1.40	0.97
1191- 5809	48.8	153.9	1.29	2.30	32.8	3.49	3.06

Table 5. Five-muon events listed by event number followed by the momenta of the outgoing muons. Odd numbered muons have the same charge as the incoming beam muon, while even numbered muons have the opposite charge. Events are produced by an incident μ^+ beam except where noted.

EVENT	P_{1X}	P_{1Y}	P_{1Z}	P_{2X}	P_{2Y}	P_{2Z}	P_{3X}	P_{3Y}	P_{3Z}	P_{4X}	P_{4Y}	P_{4Z}	P_{5X}	P_{5Y}	P_{5Z}
551- 6849	0.7	0.2	82.5	0.3	-0.2	16.6	0.1	0.0	25.4	-0.1	-0.1	8.0	0.3	0.0	23.9
623- 3285	-0.6	0.0	102.2	0.0	0.4	33.4	-0.3	-0.0	28.4	-0.0	-0.3	14.1	-0.1	-0.2	5.8
803- 6308	1.7	0.	150.1	0.4	-0.3	31.1	-0.0	-0.0	6.7	0.6	0.2	24.5	0.1	-0.0	4.5
830 9811	-2.3	0.4	137.2	-0.1	-0.2	15.5	-0.7	0.1	30.3	-0.1	-0.1	11.9	-0.0	0.3	6.3
851-11418*	-1.1	1.7	144.1	0.0	1.6	19.0	1.2	0.4	12.6	-0.1	-1.0	15.9	-0.3	-0.6	9.9
851-11970*	-2.4	-0.3	162.3	-0.0	-0.3	10.9	-0.7	0.1	30.1	0.1	0.2	5.2	-0.5	-0.4	11.9
859- 4305*	0.8	-0.1	61.8	0.4	0.8	47.4	0.5	0.4	59.7	-0.5	-0.3	36.5	-0.1	0.2	4.6
861- 206*	0.4	0.1	85.1	-1.1	1.3	38.4	-0.0	-0.1	24.0	-0.0	-0.0	12.4	0.3	1.1	7.3
890- 1460*	-0.0	-0.4	79.5	-0.5	-0.0	21.2	0.1	0.1	31.1	0.9	-0.4	20.1	-0.4	0.2	19.3
1095- 9242	1.6	0.4	106.1	0.1	0.3	25.2	0.7	0.2	25.0	0.2	0.2	22.2	0.1	-0.5	8.4

*= μ^- beam

Table 6. Five-muon events listed by event number followed by the shower energy deposited in the calorimeter, E_{shwr}, the energy transfer ν, the momentum transfer squared Q^2, the momentum of the daughter muons, $p_\perp$, perpendicular to the virtual photon direction, the missing energy, E_{miss}, and the invariant masses formed by the pairings of muons 2 and 4 with muons 3 and 5.

EVENT	E_{shwr}	ν	Q^2	$p_\perp$	E_{miss}	M_{23}	M_{25}	M_{43}	M_{45}
551- 6849	35.8	118.5	0.22	0.41	8.8	0.41	0.36	0.33	0.41
623- 3285	4.4	102.4	0.64	0.82	16.5	0.58	0.50	0.69	0.29
803- 6308	-0.1	61.0	0.02	0.39	-5.7	0.36	0.44	0.33	0.33
830- 9811	5.0	61.6	0.29	0.62	-7.4	0.59	0.39	0.63	0.51
851-11418	3.7	63.2	3.08	1.92	2.0	2.28	1.93	0.57	0.37
851-11970	9.0	45.5	0.08	0.28	-21.6	0.72	0.79	0.85	0.77
859- 4305	4.4	151.0	0.50	1.15	-4.9	1.26	1.26	0.71	0.84
861- 206	16.5	123.5	0.05	1.41	24.8	1.38	0.23	3.18	1.43
890- 1460	45.7	132.8	0.66	0.71	-4.8	0.67	1.16	0.28	1.40
1095- 9242	7.7	96.8	0.19	0.34	8.2	0.67	0.50	0.99	0.88

Table 7. Probability that the inelastic 2 muon events have different distributions in the specified kinematic variables from the inelastic odd sign 3 muon events analyzed with N_μ muons.

Variable	N_μ	Probability
Shower Energy	3	57%
ν	2 (3)	97% (97%)
Q^2	2 (3)	2% (2%)
$p_\perp$ to γ_V	2 (3)	75% (91%)
Inelasticity	2	99.6%
Missing Energy	3	42%

Table 8. Probability that the elastic 3 muon events have different distributions in the specified kinematic variables from the elastic 5 muon events analyzed with N_μ muons.

Variable	N_μ	Probability
Shower Energy	5	56%
ν	3 (5)	55% (82%)
Q^2	3 (5)	30% (38%)
$p_\perp$ to γ_V	5	½% (30%)
Inelasticity	3 (5)	58%
Missing Energy	5	63%
Invariant Mass	3 (5)	6% (15%)

Table 9. Probability that the inelastic 2 muon events have different distributions in the specified kinematic variables from the inelastic 4 muon events analyzed with N_μ muons.

Variable	N_μ	Probability
Shower Energy	4	92%
ν	4 (2)	70% (70%)
Q^2	4 (2)	37% (66%)
$p_\perp$ to γ_V	2 (4)	30% (51%)
Inelasticity	2	36%
Missing Energy	4	71%

Table 10a. Probability that the inelastic 3 muon events have different distributions in the specified kinematic variables from the inelastic 4 and 5 muon events analyzed with $\leq N_\mu$ muons.

Variable	N_μ	Probability
Shower Energy	5	99.5%
ν	5 (3)	99.9% (99.98%)
Q^2	3 (5)	82% (87%)
$p_\perp$ to γ_V	3 (5)	92% (98%)
Inelasticity	5	91%
Missing Energy	5	99%
Invariant Mass	5 (3)	66% (82%)

Table 10b. Probability that the inelastic 3 muon events have different distributions in the specified kinematic variables from the inelastic 5 muon events when they are analyzed with N_μ muons.

Variable	N_μ	Probability
Shower Energy	5	1%
ν	3 (5)	9% (40%)
Q^2	5 (3)	8% (25%)
p to γ_V	5 (3)	52% (54%)
Inelasticity	5	1%
Missing Energy	5	30%
Invariant Mass	5	60%

Table 11. Numbers of exotic multimuon events categorized by type and shower energy, E_{SH}, from data corresponding to an integrated luminosity of $0.78\times10^{39}cm^{-2}$. Also included are the expected number of events as explained in the text.

Event	E_{SH}	Number	Expected
$\mu^+\rightarrow\mu^+\mu^-\mu^\pm$	>12.5	36	39-42
$\mu^+\rightarrow\mu^+\mu^-\mu^+\mu^\pm$	<6	3	<10
	>6	13	11-27
$\mu^+\rightarrow\mu^+\mu^-\mu^+\mu^-\mu^+$	<6	5	<10
	>6	5	2

Figure Captions

Figure VII.1. Computer generated picture of odd-sign trimuon event 851-5726. Top frame: plan view; bottom frame: elevation view. Superimposed digits are the track numbers mentioned in table 1. Typically, in each interstice between modules a track registers in a proportional chamber (left tic) and, in the plan view, also in a drift chamber (right tic closest to left tic). The drift chambers are nosier due to their longer livetime. Short vertical lines at the top are calorimeter counter pulse heights. The vertical lines in the two frames are projections of trigger counters which were tagged. Heavy broken lines are tracings of the computer-reconstructed trajectories.

Figure VII.2. Computer generated picture of 4-muon event 1191-5809. Top frame: plan view; bottom frame: elevation view. Superimposed digits are the track numbers mentioned in table 3. Typically, in each interstice between modules a track registers in a proportional chamber (left tic) and, in the plan view, also in a drift chamber (right tic closest to left tic). Short vertical lines at the top are calorimeter counter pulse heights. The vertical lines in the two frames are projections of trigger counters which were tagged. Heavy broken lines are tracings of the computer-reconstructed trajectories.

Figure VII.3. Computer generated picture of 5-muon event 851-11418. Top frame: plan view; bottom frame: elevation view. Superimposed digits are the track numbers mentioned in table 5. Typically in each interstice between modules a track registers in a proportional chamber (left tic) and, in the plan view, a drift chamber (right tic closest to left tic). Tracks 3 and 4, while close in the plan view are connected

by diagonal plane wire hits to clearly separated tracks in the elevation view. Short vertical lines at the top are calorimeter counter pulse heights. The vertical lines in the two frames are projections of trigger counters which were tagged. Heavy broken lines are tracings of the computer reconstructed trajectories.

Figure VII.4. Distributions in six reconstructed kinematic variables for inelastic dimuons and the odd sign trimuons. Both types of event have shower energy greater than 12.5 GeV and energy transfer ν greater than 30 GeV. The inelastic dimuons displayed consist of all dimuons with the properly normalized π- and K-decay Monte Carlo events subtracted off. The vertical scales refer to the dimuons only. The scale for the trimuons is 2 events per division. The plain histograms represent the dimuons and the slashed columns represent the trimuons. In all distributions except (c), (d) and (e), the trimuons have had their slowest muon removed and are analyzed as dimuons. Distributions shown are (a) momentum transfer squared, (b) energy transfer ν, (c) inelasticity, (d) missing energy, (e) shower energy, and (f) the momentum of the daughter muon perpendicular to the virtual photon direction. All events pass the same standard cuts.

Figure VII.5. Feynman diagrams for the electromagnetic production of muon tridents for a target T: (a) Bethe-Heitler (b) muon bremsstrahlung, (c) target bremsstrahlung. From Ref. 15.

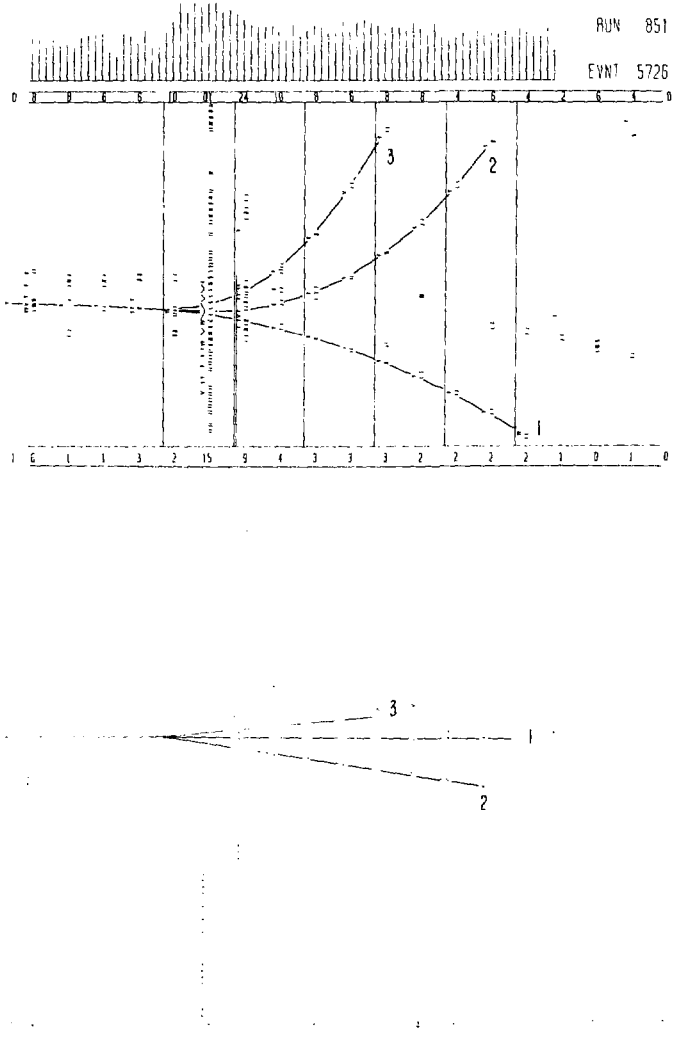
Figure VII.6. Feynman diagrams for the electromagnetic production of a muon pair in (a) an electromagnetic trident (Double Bethe-Heitler) and in (b) a charm dimuon.

Figure VII.7. Distributions in six reconstructed kinematic variables for elastic tridents and elastic 4- and 5-muon events. Both types of event have shower energies less than 6 GeV. The elastic tridents consist of all elastic trimuons with the properly normalized psi Monte Carlo events subtracted off. The vertical scales refer to the tridents only. The scale for 4- and 5-muon events is 2 events per division. The plain histograms represent the tridents and the slashed columns represent the 4- and 5-muon events. In all distributions except (c) and (d), the 4- (5-) muon events have had their slower muon(s) removed and are analyzed as tridents. Distributions shown are (a) momentum transfer squared, (b) energy transfer ν, (c) inelasticity, (d) missing energy, (e) invariant mass of the daughter muon pairs, which for the 4- and 5-muon events includes all possible pairings with the pairing produced by the two most energetic (fast) muons with the appropriate signs being shaded, and (f) the momentum of the daughter muons together perpendicular to the virtual photon direction. All events pass the same standard cuts.

Figure VII.8. Distributions in six reconstructed kinematic variables for inelastic dimuons and inelastic 4- muon events. Both types of events have shower energy greater than 12.5 GeV and energy transfer greater than 30 GeV. The inelastic dimuons displayed consist of all dimuons with the properly normalized π- and K-decay Monte Carlo events subtracted off. The vertical scales refer to the dimuons only. The scale for the 4- muon events is 2 events per division. The plain histograms represent the dimuons and the slashed columns represent the 4- muon events. In all distributions except (c), (d) and (e), the 4- muon events have had the slower muon of each sign removed and are analyzed as

dimuons. Distributions shown are (a) momentum transfer squared, (b) energy transfer ν, (c) inelasticity, (d) missing energy, (e) shower energy and (f) the momentum of the daughter muon perpendicular to the virual photon direction. All events pass the same standard cuts.

Figure VII.9. Distributions in six reconstructed kinematic variables for inelastic tridents and inelastic 4- and 5-muon events. All events have shower energy greater than 6 GeV. The inelastic tridents displayed consist of all trimuons with the properly normalized inelastic psi Monte Carlo subtracted off. The vertical scale refers to the tridents only. The scale for the 4- and 5-muon events is 2 events per division. The plain histograms represent the tridents while the left to right ascending slashed columns represent the 4-muon events and the left to right descending slashed columns represent the 5-muon events. In all distributions except (c) and (d) the 4- (5-) muon events have had their slower muon(s) removed and are analyzed as trimuons. Distributions shown are (a) momentum transfer squared, (b) energy transfer ν, (c) inelasticity, (d) missing energy, (e) invariant mass of the daughter muon pairs, which for the 4- and 5-muon events includes all possible pairings with the pairing produced by the two most energetic (fast) muons with the appropriate signs being shaded, and (f) the momentum of the daughter muons together perpendicular to the virtual photon direction. All events pass the same standard cuts.



XBL814-2255

Figure VII.1

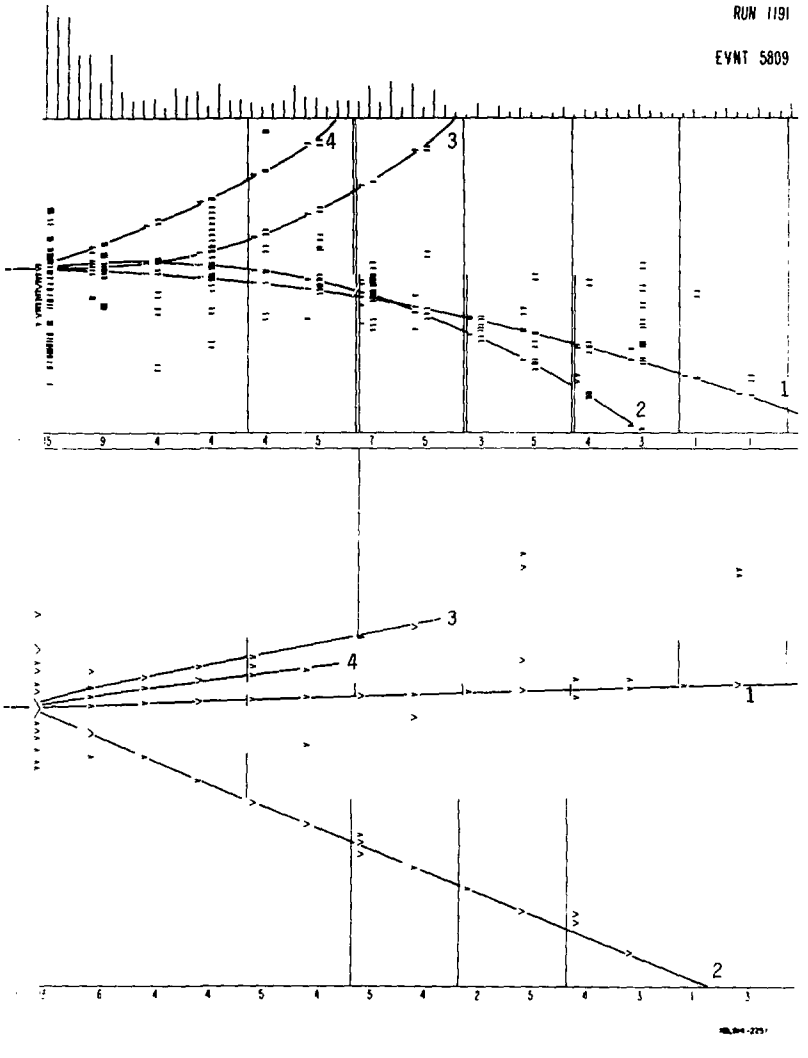
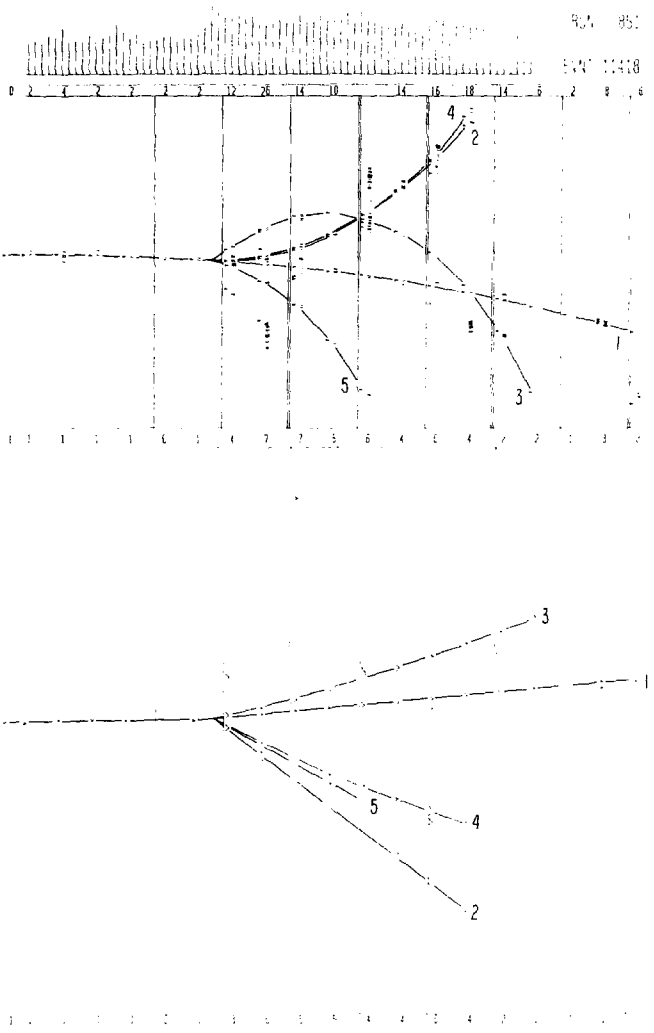


Figure VII.2



XBL814-2256

Figure VII.3

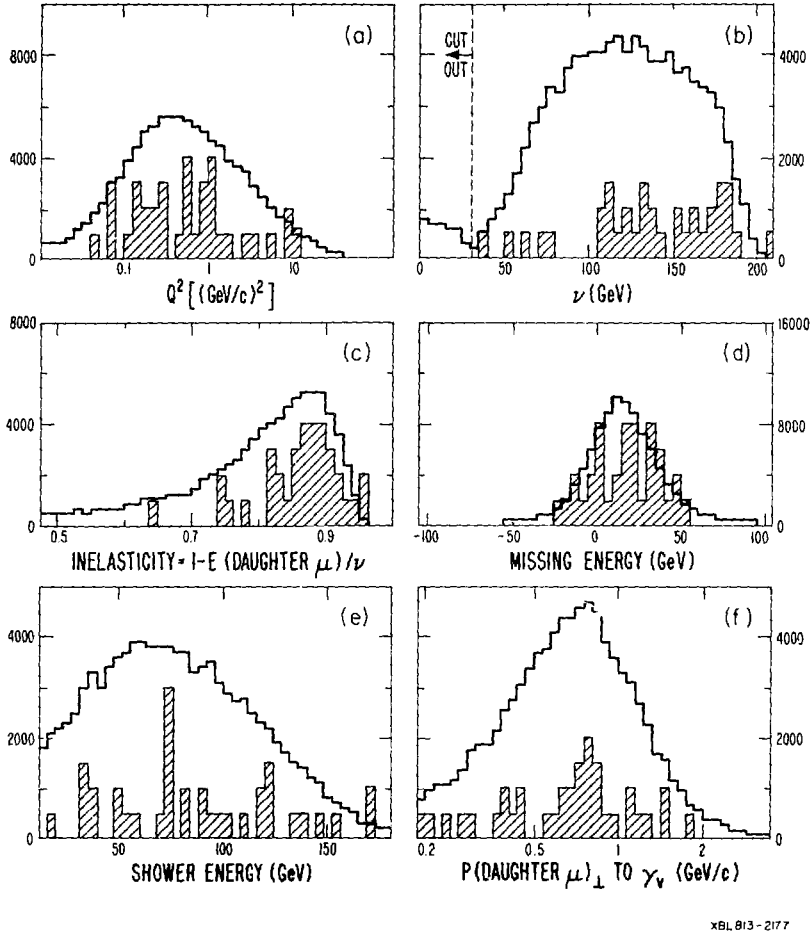
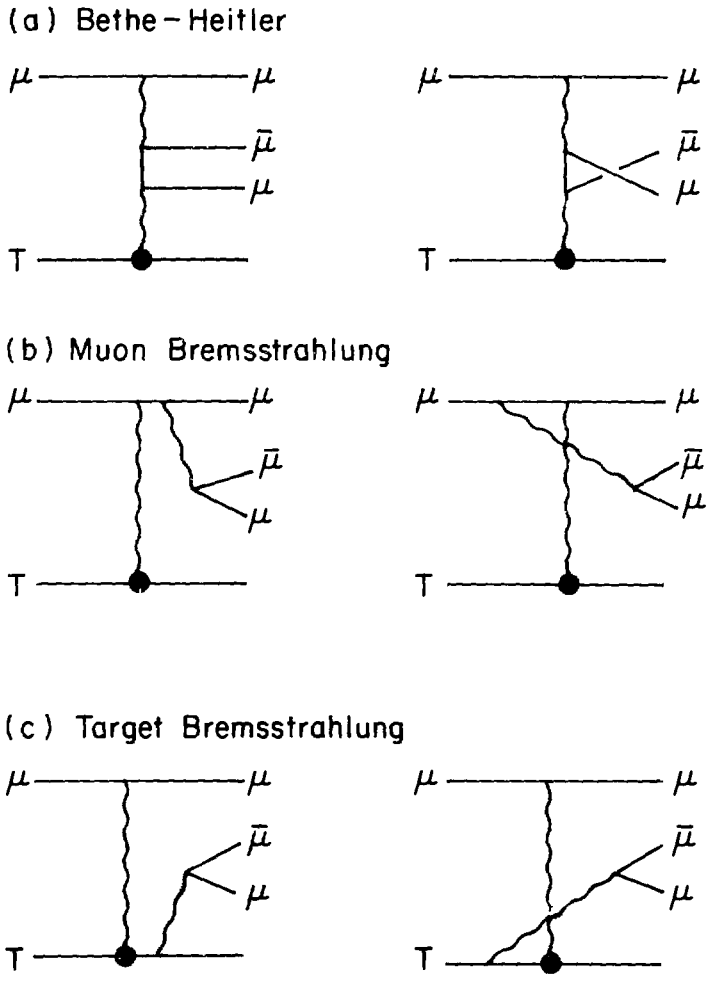

(a)
8000
4000
0
0.1
1
10
Q^2 [(GeV/c)^2]
(b)
CUT
OUT
4000
2000
0
50
100
150
200
ν (GeV)
(c)
8000
4000
0
0.5
0.7
0.9
INELASTICITY = 1-E (DAUGHTER μ)/ν
(d)
16000
8000
0
-100
-50
0
50
100
MISSING ENERGY (GeV)
(e)
4000
2000
0
50
100
150
SHOWER ENERGY (GeV)
(f)
4000
2000
0
0.2
0.5
1
2
P(DAUGHTER μ)$_\perp$ TO γ_ν (GeV/c)
XBL 813-2177

Figure VII.4



XBL814-2232

Figure VII.5

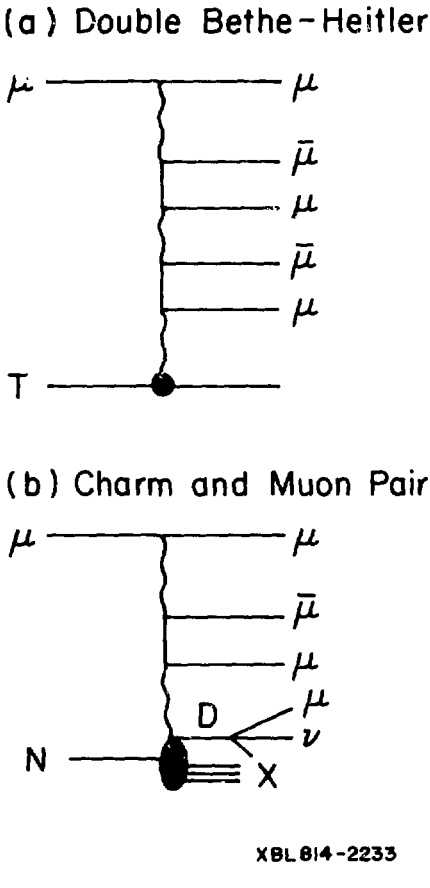

(a) Double Bethe-Heitler
μ
μ
μ̄
μ
μ̄
μ
T
(b) Charm and Muon Pair
μ
μ
μ̄
μ
D
μ
ν
N
X
XBL 814-2233

Figure VII.6

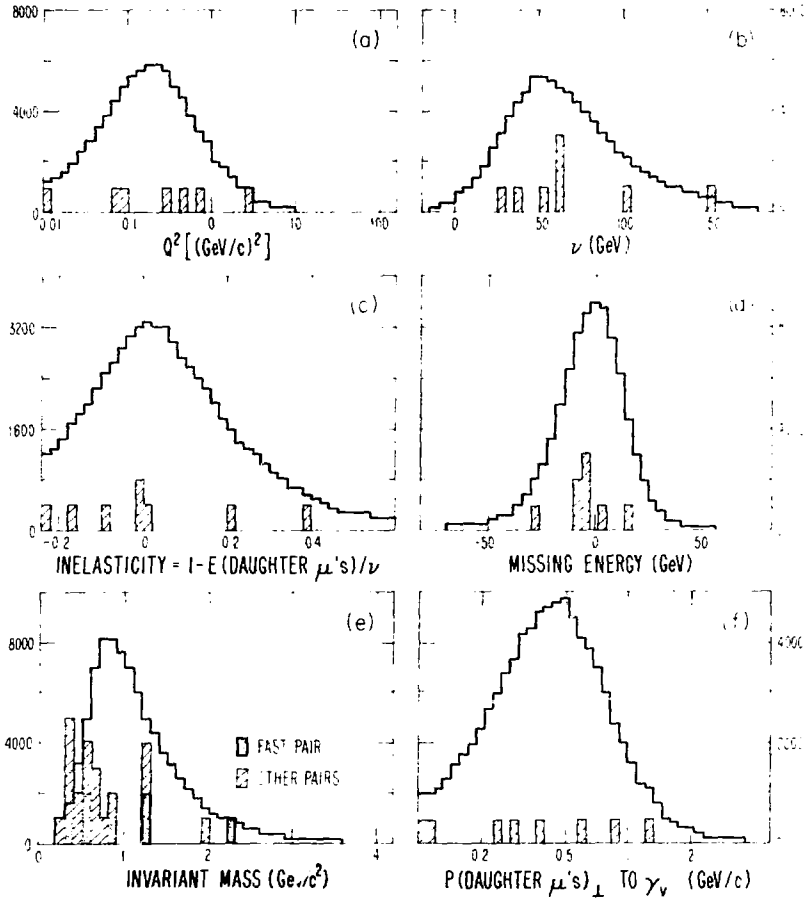

(a)
Q^2 [(GeV/c)^2]
(b)
ν (GeV)
(c)
INELASTICITY = 1 - E (DAUGHTER μ's)/ν
(d)
MISSING ENERGY (GeV)
(e)
FAST PAIR
OTHER PAIRS
INVARIANT MASS (GeV/c^2)
(f)
P (DAUGHTER μ's)$_\perp$ TO γ_ν (GeV/c)

Figure VII.7

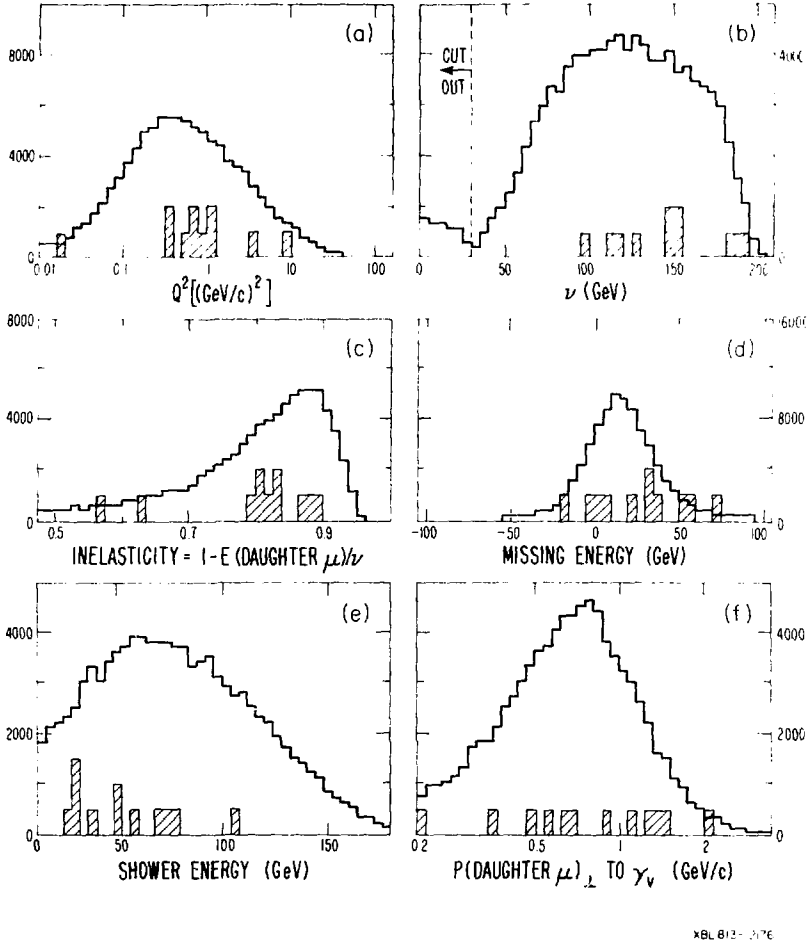

(a)
8000
4000
0
0.01
0.1
1
10
100
$Q^2[(GeV/c)^2]$
(b)
CUT
OUT
0
50
100
150
ν (GeV)
(c)
8000
4000
0
0.5
0.7
0.9
INELASTICITY = 1-E (DAUGHTER μ)/ν
(d)
8000
0
-100
-50
0
50
100
MISSING ENERGY (GeV)
(e)
4000
2000
0
0
50
100
150
SHOWER ENERGY (GeV)
(f)
4000
2000
0
0.2
0.5
1
2
P(DAUGHTER μ)$_\perp$ TO γ_ν (GeV/c)

Figure VII.8

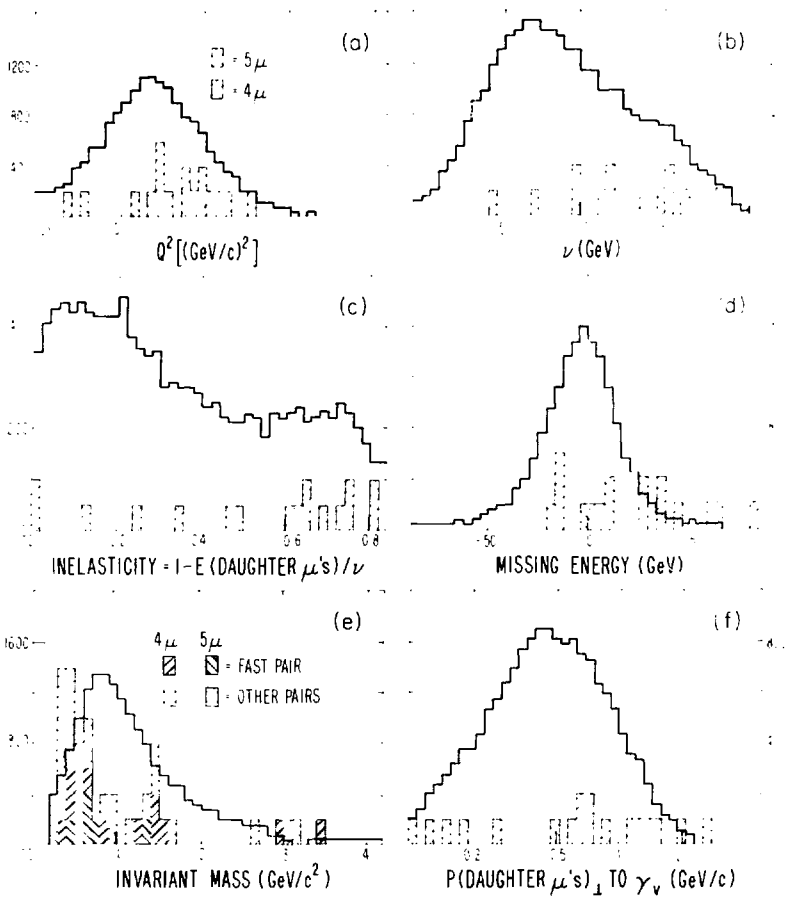

(a)
□ = 5μ
□ = 4μ
$Q^2[(GeV/c)^2]$
(b)
ν (GeV)
(c)
INELASTICITY = 1−E(DAUGHTER μ's)/ν
(d)
MISSING ENERGY (GeV)
(e)
4μ 5μ
▨ ▨ = FAST PAIR
□ = OTHER PAIRS
INVARIANT MASS (GeV/c²)
(f)
P(DAUGHTER μ's)⊥ TO γ_ν (GeV/c)

Figure VII.9